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## OSCILLATION CRITERIONS.

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1. In what follows all functions are real-valued and continuous for large positive  $t$ , say for  $t \geq t_0$ . Let  $f(t)$  be such a function. Then from STURM'S separation theorem it follows that either all non-trivial solutions of the linear differential equation

$$x'' + f(t)x = 0 \quad (1)$$

have an infinity of zeros or all have only a final number of them. In the first case (1) is called oscillatory, in the second non-oscillatory. When namely a non-trivial solution  $x(t)$  of (1) vanishes it changes its sign, since  $x(t)$  and  $x'(t)$  cannot vanish simultaneously.

It is important to know the conditions under which (1) is oscillatory. KNESER [1] already dealt with this problem in the case where  $f(t)$  has a constant sign. Recently WINTNER [2] examined the same question but without the supposition that  $f(t)$  has a constant sign for large positive  $t$ . If we denote

$$F(t) = \int f(s) ds,$$

the sufficient conditions for (1) being oscillatory are: either

$$\lim_{t \rightarrow \infty} F(t) = \infty \quad (2)$$

or only

$$\lim_{t \rightarrow \infty} \frac{\int F(s) ds}{t} = \infty. \quad (3)$$

In this note we shall deal with STURM'S differential equation

$$(Q(t) \cdot x')' + f(t)x = 0 \quad (4)$$

where  $Q(t)$  is positive, and we shall prove two oscillation criterions which are a generalization of (2). As a corollary of the first criterion we shall obtain a generalisation of a certain oscillation criterion of KNESER.<sup>1)</sup>

<sup>1)</sup> Added in proof. Two new papers containing similar results have got into my hands: J. G. MIKUSIŃSKI, On Fite's oscillation theorems, *Colloquium math.* **2** (1949) 34—39; W. LEIGHTON, The detection of the oscillation of solutions of a second order linear differential equation, *Duke Math. J.* **17** (1950) 57—61.

**2. Theorem I.** Let the integral  $\int \frac{ds}{Q(s)}$  be divergent and let exist a positive function  $\omega(t)$  having a continuous first derivative such that

$$\int \frac{Q(s)}{\omega(s)} \omega'^2(s) ds < \infty, \quad (5)$$

$$\int \omega(s) \cdot f(s) ds \rightarrow \infty \text{ for } t \rightarrow \infty. \quad (6)$$

Then (4) is oscillatory.

**Proof.** Assume (4) to be non-oscillatory. First of all we are going to prove that then there exists one unbounded solution  $x_0(t)$ .

If  $x_1(t), x_2(t)$  are two linearly independent solutions of (4), it is easy to see that their Wronskian is equal to  $\frac{W_0}{Q(t)}$ , where  $W_0$  is a non-vanishing constant, i. e.

$$x_1(t) \cdot x'_2(t) - x'_1(t) \cdot x_2(t) = \frac{W_0}{Q(t)}.$$

$x_1(t)$  and  $x_2(t)$  can be assumed such that  $W_0 = 1$ . As we suppose (4) to be non-oscillatory, the fractions  $\frac{x_1(t)}{x_2(t)}$  and  $\frac{x_2(t)}{x_1(t)}$  are continuous for large positive  $t$  and either  $\frac{x_1(t)}{x_2(t)} > 0$  or  $\frac{x_2(t)}{x_1(t)} < 0$ . In the first case it is

$$\left( \frac{x_1(t)}{x_2(t)} \right)' = - \frac{1}{Q(t) x_2^2(t)}$$

so that  $|x_2(t)|$  must be unbounded. For if  $|x_2(t)| < K$ , then

$$\begin{aligned} \frac{x_1(t)}{x_2(t)} &= \frac{x_1(t_1)}{x_2(t_1)} + \int_{t_1}^t \left( \frac{x_1(s)}{x_2(s)} \right)' ds = \frac{x_1(t_1)}{x_2(t_1)} - \int_{t_1}^t \frac{ds}{Q(s) x_2^2(s)} < \frac{x_1(t_1)}{x_2(t_1)} - \\ &\quad - \frac{1}{K^2} \int_{t_1}^t \frac{ds}{Q(s)} \rightarrow -\infty \end{aligned}$$

which is contradictory to  $\frac{x_1(t)}{x_2(t)} > 0$ . In the second case

$$\left( \frac{x_2(t)}{x_1(t)} \right)' = \frac{1}{Q(t) x_1^2(t)}$$

so that  $|x_1(t)|$  must be unbounded. Therefore at least one of the solutions  $x_1(t)$  and  $x_2(t)$  is in absolute value unbounded. We denote it  $x_0(t)$  and can assume  $x_0(t) > 0$  for large positive  $t$ , say for  $t \geq t_1$ .

Put now  $u(t) = Q(t) \frac{x_0'(t)}{x_0(t)}$ ,  $t \geq t_1$ . In view of (4)  $u(t)$  satisfies RICCATI'S differential equation

$$u' = -\frac{1}{Q(t)} u^2 - f(t). \quad (7)$$

Multiply (7) with  $\omega(t)$ , integrate the left side by parts and use SCHWARZ'S inequality and (5). One gets

$$\begin{aligned} \omega(t) \cdot u(t) &= a + \int \omega'(s) u(s) ds - \int \frac{\omega(s)}{Q(s)} u^2(s) ds - \int \omega(s) \cdot f(s) ds \\ &\leq a + \left\{ \int \frac{\omega(s)}{Q(s)} u^2(s) ds \cdot \int \frac{Q(s)}{\omega(s)} \omega'^2(s) ds \right\}^{\frac{1}{2}} - \\ &\quad - \int \frac{\omega(s)}{Q(s)} u^2(s) ds - \int \omega(s) f(s) ds \\ &\leq a + I_0^{\frac{1}{2}} \left\{ \int \frac{\omega(s)}{Q(s)} u^2(s) ds \right\}^{\frac{1}{2}} - \int \frac{\omega(s)}{Q(s)} u^2(s) ds - \int \omega(s) f(s) ds \end{aligned}$$

where  $I_0 = \int \frac{Q(s)}{\omega(s)} \omega'^2(s) ds$ , so that by (6)  $\omega(t) \cdot u(t) \rightarrow -\infty$  for  $t \rightarrow \infty$ .

Therefore  $u(t)$  and consequently  $x_0'(t)$  are negative for large positive  $t$ . But this is a contradiction because  $x_0(t)$  is positive and unbounded.

**Corollary 1.** Let  $\varepsilon$  be positive and

$$\int s^{1-\varepsilon} f(s) ds \rightarrow \infty \text{ for } t \rightarrow \infty.$$

Then (1) is oscillatory.

**Proof:** We use Theorem I with  $Q(t) = 1$  and  $\omega(t) = t^{1-\varepsilon}$ .

As a further corollary we shall derive from Theorem I a generalization of an oscillation criterion of KNESER [1]. With a slight extension, we can formulate KNESER'S result as follows:

If

$$\liminf_{t \rightarrow \infty} t^2 \cdot f(t) > \frac{1}{4}$$

then (1) is oscillatory.

Recently this criterion has been generalised simultaneously by HILLE (see [3], p. 249) and HARTMAN (see [4], p. 778). If we write

$$\log_1 t = \log t, \log_p t = \log \log_{p-1} t, p = 1, 2, 3, \dots$$

$$L_0(t) = t, L_p(t) = L_{p-1}(t) \log_p t,$$

$$S_p(t) = \sum_{n=0}^p [L_n(t)]^{-2},$$

then their sufficient condition is

$$\liminf_{t \rightarrow \infty} [L_p(t)]^2 \{f(t) - \frac{1}{4} S_{p-1}(t)\} > \frac{1}{4}.$$

Our corollary doesn't contain this result but it isn't also contained in it because it admits the possibility that the above mentioned lower limit is smaller or equal to  $\frac{1}{4}$ .

**Corollary 2.** *If*

$$\int_t^{\infty} L_p(s) \{f(s) - \frac{1}{4} S_p(s)\} ds \rightarrow \infty \text{ for } t \rightarrow \infty,$$

then (1) is oscillatory.

**Proof:** By the assumption and Theorem I with  $Q(t) = L_p(t)$  and  $\omega(t) = 1$  the differential equation

$$(L_p(t) \cdot z')' + L_p(t) \{f(t) - \frac{1}{4} S_p(t)\} \cdot z = 0 \quad (8)$$

is oscillatory. If we put  $z(t) = \frac{x(t)}{\sqrt{L_p(t)}}$ ,  $x(t)$  is also oscillatory.  $x(t)$  however satisfies (1) as we are going to prove.

From (8) it follows that  $x(t)$  satisfies the differential equation

$$x'' + \left\{ \frac{1}{4} \frac{L_p'^2(t) - 2L_p(t) \cdot L_p''(t)}{L_p^2(t)} + f(t) - \frac{1}{4} S_p(t) \right\} x = 0.$$

Therefore it is sufficient to prove that

$$\frac{L_p'^2(t) - 2L_p(t) L_p''(t)}{L_p^2(t)} = S_p(t). \quad (9)$$

For  $p = 0$  (9) is evident. Suppose that (9) holds for  $p = m - 1$ . It is

$$(\log_m t)' = \frac{1}{L_{m-1}(t)}, L'_m(t) = (L_{m-1}(t) \cdot \log_m t)' = L_{m-1}(t) \cdot \log_m t + 1,$$

$$L''_m(t) = L''_{m-1}(t) \log_m t + \frac{L'_{m-1}(t)}{L_{m-1}(t)}$$

so that

$$\begin{aligned} \frac{L_m'^2(t) - 2L_m(t) L_m''(t)}{L_m^2(t)} &= \frac{[\log_m t]^2 \cdot [L'_m(t) - 2L_{m-1}(t) \cdot L''_{m-1}(t)] + 1}{L_{m-1}^2(t) [\log_m t]^2} = \\ &= \frac{L_{m-1}'^2(t) - 2L_{m-1}(t) L_{m-1}''(t)}{L_{m-1}^2(t)} + \frac{1}{L_m^2(t)} = S_{m-1}(t) + \frac{1}{L_m^2(t)} = S_m(t). \end{aligned}$$

$$3. \limsup_{t \rightarrow \infty} \frac{\int_t^{\infty} F(s) ds}{t} = \infty \text{ is weaker than WINTNER'S supposition (3)}$$

and

$$\limsup_{t \rightarrow \infty} F(t) = \infty \quad (10)$$

is still weaker.

I have dealt with the question whether (10) suffices in order that (1) may be oscillatory. I have succeeded only in showing that (10) is a sufficient condition when  $f(t)$  is bounded from below. More is contained in the following

**Theorem II.** *Let the integral  $\int \frac{ds}{Q(s)}$  be divergent,  $Q(t) \cdot f(t)$  bounded from below and let (10) hold. Then (4) is oscillatory.*

**Proof:** Let be  $Q(t) \cdot f(t) > -M$ ,  $M > 0$ . Assume (4) to be non-oscillatory. Then by the proof of Theorem I there exists an unbounded solution  $x_0(t)$  positive from a certain  $t_1$ .

If we put  $u(t) = Q(t) \frac{x'_0(t)}{x_0(t)}$  again, we get from (7)  $u'(t) \leq -f(t)$ .

Hence a quadrature shows that

$$u(t) \leq a - F(t). \quad (11)$$

By (10) and (11) there exists  $t_2 > t_1$  such that  $u(t_2) < -\sqrt{M}$ .  $u(t)$  decreases in a certain neighbourhood of the number  $t_2$  for

$$u'(t_2) = \frac{-u^2(t_2) - Q(t_2) \cdot f(t_2)}{Q(t_2)} < \frac{-M + M}{Q(t_2)} = 0.$$

$u(t)$  however must decrease for all  $t \geq t_2$ . Let namely  $t_3$  be the first number greater than  $t_2$  and such that  $u'(t_3) = 0$ . Within the interval  $(t_2, t_3)$ ,  $u(t)$  decreases and as  $u(t_2) < -\sqrt{M}$  it is also  $u(t_3) < -\sqrt{M}$ .

Therefore

$$u'(t_3) = \frac{-u^2(t_3) - Q(t_3) \cdot f(t_3)}{Q(t_3)} < 0.$$

We have proved that  $u(t)$  decreases for  $t \geq t_2$ . Therefore

$$u(t) \leq u(t_2) < -\sqrt{M} < 0$$

for  $t \geq t_2$  and consequently  $x'_0(t) < 0$  for  $t \geq t_2$ . But this is a contradiction because  $x_0(t)$  is positive and unbounded.

**Corollary.** *Let  $f(t)$  be bounded from below and let (10) hold. Then (1) is oscillatory.*

#### References.

- [1] A. KNESER: Untersuchungen über die reellen Nullstellen der Integrale linearer Differentialgleichungen, *Math. Annalen* **42** (1893), pp. 409—435.
- [2] A. WINTNER: A criterion of oscillatory stability, *Quart. Appl. Math.* **5** (1949), pp. 115—117.
- [3] E. HILLE: Non-oscillation theorems, *Trans. Amer. Math. Soc.* **64** (1948), pp. 234—252.
- [4] P. HARTMAN: On the linear logarithmico-exponential differential equation of the second order, *Amer. J. Math.* **70** (1948), pp. 764—779.

## Oscilační kriteriia.

(Obsah předcházejícího článku.)

Nechť  $f(t)$  a  $Q(t)$  jsou reálné funkce spojité pro velká kladná  $t$  a  $Q(t) > 0$ . Pak ze STURMOVA srovnávacího teoremu plyne, že buďto všechna reálná netriviální řešení STURMOVY diferenciální rovnice (4) mají nekonečně mnoho nulových bodů nebo jen konečný jejich počet. V prvním případě nazýváme řešení oscilatorická, ve druhém případě neoscilatorická, poněvadž prochází-li řešení nulovým bodem, mění znaménko. V tomto článku jsou dokázány dvě kriteriia:

I. Nechť integrál  $\int \frac{ds}{Q(s)}$  diverguje a necht' existuje taková kladná funkce  $\omega(t)$  mající spojitou první derivaci, že

$$\int_{\infty}^{\infty} \frac{Q(s)}{\omega(s)} \omega'^2(s) ds < \infty,$$
$$\int_{\infty}^t \omega(s) f(s) ds \rightarrow \infty \text{ pro } t \rightarrow \infty.$$

Pak řešení STURMOVY diferenciální rovnice (4) jsou oscilatorická.

II. Nechť integrál  $\int \frac{ds}{Q(s)}$  diverguje, součin  $f(t) \cdot Q(t)$  je zdola ohraničen a

$$\limsup_{t \rightarrow \infty} \int f(s) ds = \infty.$$

Pak řešení Sturmovy diferenciální rovnice (4) jsou oscilatorická.

Z kriteriia I plynou dva koroláry, z nichž druhý je zobecnění jednoho KNESEROVA oscilačního kriteriia:

Korolár 1. Necht'  $\varepsilon > 0$  a

$$\int s^{1-\varepsilon} f(s) ds \rightarrow \infty \text{ pro } t \rightarrow \infty.$$

Pak řešení diferenciální rovnice (1) jsou oscilatorická.

Korolár 2. Necht'

$$\int L_p(s) \cdot \{f(s) - \frac{1}{2}S_p(s)\} ds \rightarrow \infty \text{ pro } t \rightarrow \infty.^2)$$

Pak řešení diferenciální rovnice (1) jsou oscilatorická.

<sup>2)</sup> Definici funkcí  $L_p(t)$  a  $S_p(t)$  viz v článku.