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ON THE METHOD OF THE SOLUTION OF ELASTICITY
THEORY PROBLEM FOR THE INHOMOGENEOUS AND
ANISOTROPIC BODY

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The problems of the theory of elasticity for an anisotropic and inhomogeneous body has been investigated by Sent-Venan, Foight, S.G. Lekhnitsky, S.A. Ambartsumyan, V.A. Lomakin, G.B. Coltchin and oth.

The detailed survey of the problem one can find in works [1] - [5].

In the present paper the methods of the solution of two class of problems for a common-type anisotropic and inhomogeneous body is suggested.

1. We shall consider a generalized torsion of the cylindrical prismatic homogeneous bar with rectilinear common-type anisotropy (21 or 18 elastic constants) [1]. Here is assumed, that the bar is in equilibrium on the action of a load distributed on the foreheads and torsional moment M_t . It is assumed too, that the lateral surface is free from an external force and volume forces are neglected. For the sake of simplicity the cross-section domain (Ω) assumed to be bounded and simple-connected.

Then the components of stress tensor can be founded by the formula

$$\begin{aligned} \sigma_x &= \frac{\partial^2 F}{\partial y^2}, & \sigma_y &= \frac{\partial^2 F}{\partial x^2}, & \tau_{xy} &= -\frac{\partial^2 F}{\partial x \partial y}, \\ \tau_{xz} &= \frac{\partial \psi}{\partial y}, & \tau_{yz} &= -\frac{\partial \psi}{\partial x}, \end{aligned} \quad (1.1)$$

$$\sigma_z = \frac{M_t}{a_{33}} \left(\frac{a_{34}}{I_2} x - \frac{a_{35}}{I_1} y \right) - \frac{1}{a_{33}} (a_{13} \sigma_x + \dots + a_{36} \tau_{xy}),$$

provided functions $F(x,y)$ and $\psi(x,y)$ to be known. These functions satisfy the following system of equations.

$$\begin{aligned} L_4 F + L_3 \psi &= 0 \\ L_3 F + L_2 \psi &= -2\mathcal{D} + \frac{M_t}{2a_{33}} \left(\frac{a_{34}^2}{I_2} + \frac{a_{35}^2}{I_1} \right), \end{aligned} \quad (1.2)$$

with the boundary conditions on Γ contour of Ω domain

$$\frac{\partial F}{\partial x} \Big|_r = 0, \quad \frac{\partial F}{\partial y} \Big|_r = 0, \quad \Psi \Big|_r = 0. \quad (1.3)$$

Here L_4, L_3, L_2 are elliptic-type differential operators

$$\begin{aligned} L_4 &= \beta_{22} \frac{\partial^4}{\partial x^4} - 2\beta_{26} \frac{\partial^4}{\partial x^2 \partial y^2} + (\beta_{12} + \beta_{66}) \frac{\partial^4}{\partial x^2 \partial y^2} - 2\beta_{16} \frac{\partial^4}{\partial x \partial y^3} + \beta_{11} \frac{\partial^4}{\partial y^4} \\ L_3 &= -\beta_{24} \frac{\partial^3}{\partial x^3} + (\beta_{25} + \beta_{46}) \frac{\partial^3}{\partial x^2 \partial y} - (\beta_{17} + \beta_{56}) \frac{\partial^3}{\partial x \partial y^2} + \beta_{15} \frac{\partial^3}{\partial y^3}, \\ L_2 &= \beta_{44} \frac{\partial^2}{\partial x^2} - 2\beta_{45} \frac{\partial^2}{\partial x \partial y} + \beta_{55} \frac{\partial^2}{\partial y^2}, \end{aligned} \quad (1.4)$$

where β_{ij} - elastic constants satisfy well-known inequality [4]

$$\beta_{ii} > 0, \quad \beta_{ii} \beta_{jj} - \beta_{ij}^2 > 0, \quad (i, j = 1, 2, 4, 5, 6). \quad (1.5)$$

Further we assume that

$$\begin{aligned} \beta_{14} &= \xi \beta_1, \quad \beta_{24} = \xi \beta_2, \quad \beta_{34} = \xi \beta_3, \quad \beta_{46} = \xi \beta_4, \quad \beta_{15} = \xi \beta_5, \\ \beta_{25} &= \xi \beta_6, \quad \beta_{35} = \xi \beta_7, \quad \beta_{56} = \xi \beta_8, \quad (\xi < 1), \end{aligned} \quad (1.6)$$

where ξ is a small parameter and $\xi = 0$ if the plane of cross-section is a plane of elastic symmetry.

Then the system of differential equations (1.2) may be written as

$$\begin{aligned} L_4 F + \xi \tilde{L}_3 \Psi &= 0 \\ \xi \tilde{L}_3 F + L_2 \Psi &= -2\vartheta + \xi^2 \tilde{A}, \end{aligned} \quad (1.7)$$

where following notations are used

$$\begin{aligned} \tilde{L}_3 &= -\beta_2 \frac{\partial^3}{\partial x^3} + (\beta_6 + \beta_4) \frac{\partial^3}{\partial x^2 \partial y} - (\beta_1 + \beta_8) \frac{\partial^3}{\partial x \partial y^2} + \beta_5 \frac{\partial^3}{\partial y^3}, \\ \tilde{A} &= \frac{M_t}{2} \left(\frac{\beta_3}{I_2} + \frac{\beta_{10}}{I_1} \right), \quad a_{34} = \xi \beta_9 \sqrt{a_{33}}, \quad a_{35} = \xi \beta_{10} \sqrt{a_{33}}. \end{aligned} \quad (1.8)$$

The solution of (1.7) is expanded in power series in small physical parameter ξ as follows

$$\Psi = \Psi_0 + \xi \Psi_1 + \xi^2 \Psi_2 + \dots \quad (1.9)$$

$$F = F_0 + \xi F_1 + \xi^2 F_2 + \dots \quad (1.10)$$

Then for the determination of $F_k(x, y), \Psi_k(x, y)$ functions from equation (1.3) and (1.7) one can obtain the follow system of recurrent boundary problems

$$\begin{aligned} L_4 F_0 &= 0, \quad (x, y) \in \Omega, \\ \frac{\partial F_0}{\partial x} &= \frac{\partial F_0}{\partial y} = 0, \quad (x, y) \in \Gamma; \end{aligned} \quad (1.11)$$

$$\begin{aligned} L_4 F_k &= f_k(x, y), \quad (f_k = -\tilde{L}_3 \Psi_{k-1}), \quad (x, y) \in \Omega, \\ \frac{\partial F_k}{\partial x} &= \frac{\partial F_k}{\partial y} = 0, \quad (k=1, 2, 3, \dots), \quad (x, y) \in \Gamma; \end{aligned} \quad (1.12)$$

$$\begin{aligned} L_2 \Psi_k &= g_k(x, y), \quad (x, y) \in \Omega, \\ \Psi_k &= 0, \quad (x, y) \in \Gamma, \end{aligned} \quad (1.13)$$

(k=1, 2, 3...)

where

$$g_0 = -2\vartheta, \quad g_2 = -L_3 F_1 + \tilde{A}, \quad g_k(x, y) = -L_3 F_{k-1}, \quad (k=1, 2, 3, \dots). \quad (1.14)$$

It follows from (1.11) that $F_0 = 0$. Thus, the solution of the problem of a generalized torsion of a homogeneous bar with a very common-type anisotropy reduce to a number of problems similar to problems of bars torsion (1.13) and plane problem for a body with special-type anisotropy (with one plane of symmetry at each point—13 or 12 elastic constants).

In view of Sobolev's theorem [6] and some results of works [4, 7] following theorems can be proved.

Theorem 1. There exists a solution of the boundary problem (1.13) of the form (1.9). If $\xi < \xi_0$, the series (1.9) converge absolutely and uniformly in x, y on Ω domain, the first derivative of (1.9) in x, y converge in every $L_p(\Omega)$, ($p > 1$) and the second derivative - in $L_2(\Omega)$ only.

Theorem 2. If $\xi < \xi_0$, then the solution of the problem (1.12) have the form (1.10). The first and second derivatives of this solution converge uniformly in Ω domain by x, y , the third derivative - in any $L_p(\Omega)$ ($p > 1$), and the fourth - in $L_2(\Omega)$ only.

Here ξ_0 is some constant depending on coefficients and constants entering in Sobolev's theorems.

2. We consider now a generalized torsion of anisotropic bar with a weak inhomogeneity [4, 7]

$$\beta_{ij} = \beta_{ij}^0 [\alpha_{ij} + \delta f_{ij}(x, y)], \quad (2.1)$$

where β_{ij}^0 are constants, $f_{ij}(x, y)$ - continuous known functions of x, y with second order derivatives at least, δ - some small physical parameter. Then the solution of the problem we seek as a power series in δ as follows.

$$F = F_0 + \delta F_1 + \delta^2 F_2 + \dots$$

$$\Psi = \Psi_0 + \delta \Psi_1 + \delta^2 \Psi_2 + \dots$$

In order to determinate F_k, Ψ_k functions we obtain the system of recurrent boundary problems, corresponding to homogeneous ones. These problems can be solved using the methods mentioned above.

Finally we note, there are convergence theorems for F, Ψ series on the analogy of 1 section.

References

1. Лехницкий С.Г. Теория упругости анизотропного тела. Изд-во "Наука", М., 1977.
2. Амбарцумян С.А. Общая теория анизотропных оболочек. Изд-во "Наука", М., 1974.
3. Ломакин В.А. Теория упругости неоднородных тел. Изд-во МГУ, М., 1976.
4. Саркисян В.С. Некоторые задачи математической теории упругости анизотропного тела. Изд-во ЕГУ, Е., 1976.
5. Теория упругости неоднородных тел. Библиографический указатель отечественной и зарубежной литературы (Составители Г.Б. Колчин Э.А. Фаверман. Изд-во "Штинца", Кишинев, 1972).
6. Соболев С.Л. Некоторые применения функционального анализа к математической физике. Изд-во ЛГУ, Л., 1950.
7. V.S. Sarkisian a A. Mamrillova "VYBRANE" KAPITOLYZ TEORIE PRUZNOSTI ANIZOTROPNYCH TELIES Universita Komenskoho v Bratislava 1979.