

Sylvia Pulmannová

On the products of quantum logics

In: Zdeněk Frolík (ed.): Proceedings of the 11th Winter School on Abstract Analysis. Circolo Matematico di Palermo, Palermo, 1984. Rendiconti del Circolo Matematico di Palermo, Serie II, Supplemento No. 3. pp. [231]--235.

Persistent URL: <http://dml.cz/dmlcz/701316>

Terms of use:

© Circolo Matematico di Palermo, 1984

Institute of Mathematics of the Academy of Sciences of the Czech Republic provides access to digitized documents strictly for personal use. Each copy of any part of this document must contain these *Terms of use*.



This paper has been digitized, optimized for electronic delivery and stamped with digital signature within the project *DML-CZ: The Czech Digital Mathematics Library* <http://project.dml.cz>

ON THE PRODUCTS OF QUANTUM LOGICS

SYLVIA PULMANOVÁ

A definition of a product of quantum logics is formulated and a comparison with the free orthodistributive product of orthomodular σ -lattices is given.

A quantum logic is the couple (L, M) , where L is an orthomodular σ -lattice and M is a set of states (i.e. probability measures) on L , which is strong for L , i.e. the statement

$$\{m \in M : m(a) = 1\} \subset \{m \in M : m(b) = 1\}$$

implies that $a \leq b$, $a, b \in L$. This notion was introduced by Gudder [4]. The physical interpretation of (L, M) is as follows. The set L is interpreted as the set of all experimentally verifiable propositions of a physical system (the "logic" of the system), and M is the set of physical states. The requirement of the existence of a strong set of states restricts the choice of orthomodular σ -lattices suited for description of physical systems, there are orthomodular σ -lattices with no states [3]. We shall also suppose the Jauch - Piron property, i.e.

$$m(a_i) = 1 \text{ for all } i \in N \text{ iff } m(\bigwedge_{i \in N} a_i) = 1$$

for any $m \in M$. More details on quantum logic can be found in [6] and [9].

To describe a physical system which is composed of two other systems, we need a kind of the product of quantum logics. In the traditional approach to quantum theory, it is supposed that to any physical system there is a Hilbert space (complex, separable, with the dimension at least three). The set of propositions is the lattice $L(H)$ of all closed linear subspaces of H and states are represented by the density operators. The joint physical system consisting of two other systems is then described by the tensor product of the Hilbert spaces of these two systems.

A product of orthomodular σ -lattices was defined in the following way [5]. We recall that the elements a, b of an orthomodular lattice L are compatible ($a \leftrightarrow b$) if

$$a = (a \wedge b) \vee (a \wedge b^\perp) .$$

Definition 1 . Let \mathcal{C} be a subcategory of the category of orthomodular σ -lattices. Let $\{L_i : i \in I\}$ and L be elements of \mathcal{C} . Then $(L, (u_i)_{i \in I})$ is a tensor product (or free orthodistributive product) of the L_i 's if

(i) $u_i : L_i \rightarrow L$ are injections in \mathcal{C} , $i \in I$,

(ii) $\bigcup_{i \in I} u_i(L_i)$ generates L ,

(iii) for any at most countable subset F of I ,

$\bigwedge_{i \in F} u_i(a_i) = 0$ for $a_i \in L_i$ iff at least one a_i is zero,

(iv) $u_i(a_i) \leftrightarrow u_j(a_j)$ for any $i, j \in I$, $i \neq j$.

In the category of Hilbert space lattices, it was shown [1] , [5] that: if H_1 , H_2 are complex, separable, of the dimension at least three, then there are exactly two (unequivalent) products, defined by

$$\begin{aligned} \text{(i)} \quad u_1 : L(H_1) &\rightarrow L(H_1 \otimes H_2) \\ P &\mapsto P \otimes H_2 \\ u_2 : L(H_2) &\rightarrow L(H_1 \otimes H_2) \\ P &\mapsto H_1 \otimes P \end{aligned}$$

and

$$\begin{aligned} \text{(ii)} \quad u_1 : L(H_1) &\rightarrow L(\overline{H_1} \otimes H_2) \\ P &\mapsto \overline{P} \otimes H_2 \\ u_2 : L(H_2) &\rightarrow L(\overline{H_1} \otimes H_2) \\ P &\mapsto \overline{H_1} \otimes P \end{aligned}$$

where $H_1 \otimes H_2$ is the tensor product of Hilbert spaces H_1 and H_2 , and \overline{H} is the dual of H .

In the case of real Hilbert spaces there is exactly one product defined by (i) .

We shall introduce a definition of a product of quantum logics. We need some preliminary remarks. Let S be a set of states on a logic L . We say that a state p on L is a superposition of the states in S if $S(a) = 1$ implies $p(a) = 1$ for $a \in L$, where $S(a) = 1$ means that $s(a) = 1$ for any $s \in S$ [9] . If (L, M) is a quantum logic, we shall write $\overline{S} = \{p \in M : S(a) = 1 \Rightarrow p(a) = 1\}$ for any $S \subset M$.

Definition 2 . Let (L_1, M_1) , (L_2, M_2) , (L, M) be quantum logics. We shall say that $(L, M)_{\alpha, \beta}$ is the tensor product of (L_i, M_i) , i

if

$$(i) \quad \alpha : L_1 \times L_2 \rightarrow L, \quad \beta : M_1 \times M_2 \rightarrow M$$

$$\beta(m_1, m_2)(\alpha(a_1, a_2)) = m_1(a_1) \cdot m_2(a_2)$$

for any $m_i \in M_i, a_i \in L_i, i=1,2,$

$$(ii) \quad \{m \in M : m(a) = 1\} = \{\beta(m_1, m_2) : \beta(m_1, m_2)(a) = 1\}^-$$

for elements $a \in L$ of the form

$$a = \bigwedge_k \alpha(a_1^k, a_2^k), \quad a_1^k \in L_1, \quad a_2^k \in L_2, \quad k \in \mathbb{N},$$

and

$$a = \alpha(a_1, 1)^{\perp}, \text{ resp. } a = \alpha(1, a_2)^{\perp}, \quad a_1 \in L, \quad a_2 \in L_2,$$

$$(iii) \quad \alpha[L_1 \times L_2] \text{ generates } L,$$

$$(iv) \quad \beta[M_1 \times M_2]^- = M.$$

Theorem 1. Let $(L, M)_{\alpha, \beta}$ be a tensor product of (L_1, M_1) and (L_2, M_2) . Let us put

$$u_1 : L_1 \rightarrow L, \quad u_2 : L_2 \rightarrow L$$

$$a \mapsto \alpha(a, 1) \quad a \mapsto \alpha(1, a).$$

Then (L, u_1, u_2) is a free product of L_1, L_2 by Def 1.

Proof. By (i) of Def. 2 we have

$$\beta(m_1, m_2)(\alpha(1, 1)) = m_1(1) \cdot m_2(1) = 1$$

for any $m_i \in M_i, i=1,2$. From this we get $m(\alpha(1, 1)) = 1$ for all $m \in \beta[M_1 \times M_2]^-$, and by (iv), $m(\alpha(1, 1)) = 1$ for all $m \in M$, i.e. $\alpha(1, 1) = u_1(1) = u_2(1) = 1$.

For any $a \in L_1, \beta(m_1, m_2)(\alpha(a^{\perp}, 1)) = m_1(a^{\perp}) \cdot m_2(1) = 1 = (1 - m_1(a)) \cdot m_2(1) = 1 - m_1(a) \cdot m_2(1) = 1 - \beta(m_1, m_2)(\alpha(a, 1)) = \beta(m_1, m_2)(\alpha(a, 1)^{\perp})$ for all $m_i \in M_i, i=1,2$. From this we obtain

$$\{\beta(m_1, m_2) : \beta(m_1, m_2)(\alpha(a^{\perp}, 1)) = 1\}^- =$$

$$\{\beta(m_1, m_2) : \beta(m_1, m_2)(\alpha(a, 1)^{\perp}) = 1\}^-,$$

which implies by (ii) that $\alpha(a^{\perp}, 1) = \alpha(a, 1)^{\perp}$, i.e.

$$u_1(a^{\perp}) = u_1(a)^{\perp}. \text{ Similarly, } u_2(a^{\perp}) = u_2(a)^{\perp}, \quad a_2 \in L_2.$$

By the Jauch - Piron property we have

$$\beta(m_1, m_2)(\alpha(\bigwedge_k a_1^k, 1)) = 1 \text{ iff } m_1(\bigwedge_k a_1^k) = 1 \text{ iff}$$

$$\beta(m_1, m_2)(\bigwedge_k \alpha(a_1^k, 1)) = 1. \text{ From this we obtain}$$

$$\{\beta(m_1, m_2) : \beta(m_1, m_2)(\alpha(\bigwedge_k a_1^k, 1)) = 1\}^- =$$

$$= \{\beta(m_1, m_2) : \beta(m_1, m_2)(\bigwedge_k (\alpha(a_1^k, 1))) = 1\}^-,$$

which implies by (ii) that $\alpha(\bigwedge_k a_1^k, 1) = \bigwedge_k \alpha(a_1^k, 1)$, i.e.

$$u_1(\bigwedge_k a_1^k) = \bigwedge_k u_1(a_1^k). \text{ This shows that } u_1 \text{ and } u_2 \text{ are orthohomomorphisms.}$$

Now $u_1(a) = u_1(a')$, $a', a \in L_1$ implies that

$$\beta(m_1, m_2)(\alpha(a, 1)) = \beta(m_1, m_2)(\alpha(a', 1)) \text{ for any}$$

$$m_i \in M_i, i=1,2, \text{ which implies that } m_1(a) = m_1(a') \text{ for any}$$

$m_1 \in M_1$, i.e. $a = a'$. Hence u_1 and u_2 are injections.

For any $m_i \in M_i, i=1,2$, we have by the Jauch - Piron property,

$\beta(m_1, m_2)(u_1(a_1) \wedge u_2(a_2)) = \beta(m_1, m_2)(\alpha(a_1, 1) \wedge \alpha(1, a_2)) = 1$
 iff $\beta(m_1, m_2)(u_1(a_1)) = 1$ and $\beta(m_1, m_2)(u_2(a_2)) = 1$ iff
 $m_1(a_1) = 1$ and $m_2(a_2) = 1$ iff $\beta(m_1, m_2)(\alpha(a_1, a_2)) = 1$, i.e.

$$\left\{ \beta(m_1, m_2) : \beta(m_1, m_2)(\alpha(a_1, 1) \wedge \alpha(1, a_2)) = 1 \right\} = \\ = \left\{ \beta(m_1, m_2) : \beta(m_1, m_2)(\alpha(a_1, a_2)) = 1 \right\},$$

hence by (ii) $u_1(a_1) \wedge u_2(a_2) = \alpha(a_1, a_2)$. This shows (ii) of Def. 1. Now let $a_1 \in L_1, a_1 \neq 0$ and $u_1(a_1) \wedge u_2(a_2) = 0, a_2 \in L_2$. Let $m_1^0 \in M_1$ be such that $m_1^0(a_1) = 1$ (the existence of m_1^0 follows from the fact that M_1 is strong for L_1). Then

$$\beta(m_1^0, m_2)(u_1(a_1) \wedge u_2(a_2)) = \beta(m_1^0, m_2)(\alpha(a_1, a_2)) = \\ = m_1^0(a_1)m_2(a_2) = 0 \text{ iff } m_2(a_2) = 0. \text{ Thus } u_1(a_1) \wedge u_2(a_2) = 0 \text{ implies } \\ m_2(a_2) = 0 \text{ for any } m_2 \in M, \text{ i.e. } a_2 = 0.$$

Finally, for any $a_i \in L_i, i=1,2$, we have

$\beta(m_1, m_2)(u_1(a_1) \wedge u_2(a_2)) = \beta(m_1, m_2)(\alpha(a_1, a_2)) = \\ = m_1(a_1)m_2(a_2) = \beta(m_1, m_2)(u_1(a_1))\beta(m_1, m_2)(u_2(a_2))$ for any $m_i \in M_i, i=1,2$. This implies that $u_1(a_1)$ and $u_2(a_2)$ are independent (in the probabilistic sense), and by [2] they have joint probability distributions in all states of $\beta[M_1 \times M_2] = M$, hence they are compatible. This completes the proof.

Let H be a real or complex separable Hilbert space, $\dim H \geq 3$. If we put $M = \{m_\varphi : \varphi \in H, \|\varphi\| = 1\}$, where m_φ is the vector state corresponding to the vector φ by the Gleason theorem [9], then $(L(H), M)$ is a quantum logic. Let $(L(H_1), M_1), (L(H_2), M_2), (L(H_1 \otimes H_2), M)$ and $(L(\overline{H}_1 \otimes H_2), \overline{M})$ be quantum logics of the corresponding Hilbert spaces. If we put

$$(i) \alpha : (P_1, P_2) \rightarrow P_1 \otimes P_2, P_i \in L(H_i), i=1,2, \\ \beta : (m_{\varphi_1}, m_{\varphi_2}) \mapsto m_{\varphi_1 \otimes \varphi_2}, \varphi_i \in H_i, i=1,2,$$

or

$$(ii) \overline{\alpha} : (P_1, P_2) \rightarrow \overline{P}_1 \otimes P_2, P_i \in L(H_i), i=1,2, \\ \overline{\beta} : (m_{\varphi_1}, m_{\varphi_2}) \mapsto m_{\overline{\varphi}_1 \otimes \varphi_2}, \varphi_i \in H_i, i=1,2,$$

then it can be easily checked that $(L(H_1 \otimes H_2), M)_{\alpha, \beta}$ and $(L(\overline{H}_1 \otimes H_2), \overline{M})_{\overline{\alpha}, \overline{\beta}}$ are the products of $(L(H_i), M_i), i=1,2$.

More details on the products of quantum logics are in [7] and [8].

REFERENCES

- [1] D. Aerts - I. Daubechies : Physical justification for using the tensor product, *Helv. Phys. Acta* 51 (1978) 661-675.
 [2] A. Dvurečenskij - S. Pulmannová : On joint distributions of obser-

vables, Math. Slovaca 32 (1982) 155-166.

[3] R. Greechie : Orthomodular lattices admitting no states, Journ.Comb.Theory 10 (1971) 119-132.

[4] S.P. Gudder : A superposition principle in physics, J.Math.Phys. 11 (1970) 1037-1040.

[5] T. Matolcsi : Tensor product of Hilbert lattices and free orthodistributive product of orthomodular lattices, Acta Sci.Math. 37 (1975) 263-272.

[6] C. Piron : Foundations of Quantum Physics.W.A. Benjamin Inc., Reading, Mass. 1976.

[7] S.Pulmannová : On the coupling of quantum logics, Int. J. Theor. Phys., in print.

[8] S. Pulmannová : Tensor products of quantum logics, in preparation.

[9] V.S. Varadarajan : Geometry of Quantum Theory, Van Nostrand, Princeton, N.J. 1968.

Mathematical Institute
Slovak Academy of Sciences
Obrancov mieru 49
814 73 Bratislava
ČSSR