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Infinitesimal Bending of a Subspace of a Space with Non-Symmetric Basic Tensor

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Abstract

In this work infinitesimal bending of a subspace of a generalized Riemannian space (with non-symmetric basic tensor) are studied. Based on non-symmetry of the connection, it is possible to define four kinds of covariant derivative of a tensor. We have obtained derivation formulas of the infinitesimal bending field and integrability conditions of these formulas (equations).

Key words: Generalized Riemannian space, infinitesimal bending, infinitesimal deformation, subspace.

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0 Introduction

0.1. A generalized Riemannian space GR_N is a differentiable manifold, endowed with non-symmetric basic tensor $G_{ij}(x^1, \dots, x^N)$ [2], whose symmetric part is $G_{\underline{ij}}$, and antisymmetric part $G_{\underset{\vee}{ij}}$.

By equations

$$x^i = x^i(u^1, \dots, u^M) \equiv x^i(u^\alpha), \quad \text{rank}(B_\alpha^i) = M, \quad (B_\alpha^i = \partial x^i / \partial u^\alpha), \quad (0.1)$$

in local coordinates is defined a *subspace* $GR_M \subset GR_N$, with metric tensor

$$g_{\alpha\beta} = B_\alpha^i B_\beta^j G_{ij}, \quad (0.2)$$

which is generally also non-symmetric. Remark that in the present work Latin indices i, j, k, \dots take values $1, \dots, N$, while Greek indices $\alpha, \beta, \gamma, \dots$ take values $1, \dots, M$, ($M < N$) and refer to the subspace.

For the lowering and raising of indices in GR_N one uses the tensor G_{ij} respectively G^{ij} , where $(G^{ij}) = (G_{ij})^{-1}$.

Christoffel symbols at GR_N are

$$\Gamma_{i.jk} = \frac{1}{2}(G_{ji,k} - G_{jk,i} + G_{ik,j}), \quad \Gamma_{jk}^i = G^{ip}\Gamma_{p.jk}, \quad (0.3a, b)$$

where, by the comma a partial derivative is denoted.

The scalar product and the orthogonality one expresses in usual way in the GR_N by G_{ij} , and in the GR_M by $g_{\alpha\beta}$.

On subspaces of generalized Riemannian spaces there exist many works, eg. [7]–[16], [19]–[23]. The present work is continuation and widening of our work [21].

0.2. If in the points of GR_M a vector field $z^i(u^\alpha)$ is defined, the equations

$$\bar{x}^i = x^i(u^\alpha) + \varepsilon z^i(u^\alpha), \quad (0.4)$$

where ε is an infinitesimal, define an *infinitesimal deformation* of the subspace GR_M . Obtained subspace will be denoted \overline{GR}_M . The vector field $z^i(u^\alpha)$ is an *infinitesimal deformation field*. In this study of infinitesimal deformations, according to (0.4), magnitudes of a degree higher than the first with respect to ε are omitted.

Among numerous, we refer on papers on infinitesimal deformations of spaces and subspaces, and related topics [4]–[9], [17], [18], [21]–[23].

0.3. A particular case of infinitesimal deformations is *infinitesimal bending* (see e.g. [7], [8], [9], [21]). By virtue of (0.4), for $\bar{g}_{\alpha\beta}$ one obtains [21]:

$$\bar{g}_{\alpha\beta} = g_{\alpha\beta} + \varepsilon(B_\alpha^i B_\beta^j G_{ij,k} z^k + B_\alpha^i z_\beta^j G_{ij} + z_\alpha^i B_\beta^j G_{ij}) \quad (0.5)$$

and, by definition, the subspace $\overline{GR}_M \subset GR_N$ is *infinitesimal bending of the subspace* $GR_M \subset GR_N$ iff (the equation (1.5) in [21]):

$$G_{ij,k} z^k B_\alpha^i B_\beta^j + G_{ij}(B_\alpha^i z_\beta^j + z_\alpha^i B_\beta^j) = 0, \quad (0.6)$$

1 Derivational formulas of the bending field

1.0. Let be $GR_M \subset GR_N$, where GR_M is defined by virtue of (0.1). Consider at points of GR_M $N - M$ mutually orthogonal unit vectors N_A^i , ($A = M + 1, \dots, N$), which are also orthogonal to GR_M , i.e. to the vectors $B_\alpha^i = \partial x^i / \partial u^\alpha$. So, here we are using also the third kind of indices:

$$A, B, C \dots \in \{M + 1, \dots, N\}.$$

From the exposed, we have the relations

$$G_{\underline{ip}}G^{pj} = \delta_i^j, \quad g_{\underline{\alpha\pi}}g^{\pi\beta} = \delta_\alpha^\beta, \quad (1.1a, b)$$

$$G_{\underline{ij}}N_A^iB_\alpha^j = 0, \quad G_{\underline{ij}}N_A^iN_B^j = e_A\delta_{AB}, \quad (e_A = \pm 1), \quad (1.2a, b)$$

where $g^{\underline{\alpha\beta}}$ is obtained analogously to G^{ij} . Similarly to (0.3), we can define Cristoffel symbols $\tilde{\Gamma}_{\beta\gamma}^\alpha$ by means of $g_{\alpha\beta}$. These symbols are in general also non-symmetric. Based on that, for a tensor defined in the points of the subspace we have 4 kinds of covariant derivative. For example [13]:

$$B_{\alpha|\mu}^i = B_{\alpha,\mu}^i + \Gamma_{pm}^i B_\alpha^p B_\mu^m - \tilde{\Gamma}_{\alpha\mu}^{\pi} B_\pi^i \quad (1.3a-d)$$

$\begin{matrix} 1 \\ 2 \\ 3 \\ 4 \end{matrix}$

$\begin{matrix} mp \\ pm \\ mp \end{matrix}$

$\begin{matrix} \mu\alpha \\ \mu\alpha \\ \alpha\mu \end{matrix}$

$$N_A^i|_\mu = N_A^i|_3\mu = N_{A,\mu}^i + \Gamma_{mp}^i N_A^p B_\mu^m. \quad (1.4a, b)$$

$\begin{matrix} 1 \\ 2 \end{matrix}$

$\begin{matrix} 3 \\ 4 \end{matrix}$

From here one obtains 4 kinds of *derivational formulae* of the subspace $GR_M \subset GR_N$ [13,14]:

$$B_{\alpha|\theta}^i = \Phi_{\theta\alpha\mu}^\pi B_\pi^i + \sum_{A=M+1}^N \Omega_{A\alpha\mu} N_A^i, \quad (1.5a)$$

$$N_{B|\theta}^i = -e_B g^{\pi\sigma} \Omega_{B\sigma\mu} B_\pi^i + \sum_{A=M+1}^N \Psi_{\theta AB\mu} N_A^i, \quad \Psi_{\theta BB\mu} = 0, \quad (1.5b)$$

where $\theta \in \{1, 2, 3, 4\}$ designates the kind of covariant derivative. With respect to (4a,b) is:

$$\Omega_{A\alpha\beta}^1 = \Omega_{A\alpha\beta}^3 \quad (1.6a, b)$$

$\begin{matrix} 1 \\ 2 \end{matrix}$

$\begin{matrix} 3 \\ 4 \end{matrix}$

$$\Psi_{AB\mu}^1 = \Psi_{AB\mu}^3 \quad (1.7a, b)$$

$\begin{matrix} 1 \\ 2 \end{matrix}$

$\begin{matrix} 3 \\ 4 \end{matrix}$

and by virtue of (48') in [13]:

$$\begin{aligned} \Phi_{2\beta\gamma}^\alpha &= -\Phi_{1\beta\gamma}^\alpha, & \Phi_{3\beta\gamma}^\alpha &= \Phi_{1\beta\gamma}^\alpha + 2\tilde{\Gamma}_{\beta\gamma}^\alpha, \\ \Phi_{4\beta\gamma}^\alpha &= -\Phi_{1\beta\gamma}^\alpha - 2\tilde{\Gamma}_{\beta\gamma}^\alpha \end{aligned} \quad (1.8 a,c)$$

1.1. The infinitesimal bending field z^i can be expressed by tangential and normal component with respect to GR_M :

$$z^i = p^\sigma B_\sigma^i + \sum_A q_A N_A^i. \quad (1.9)$$

Using this value, the condition (0.6) becomes

$$\begin{aligned}
& G_{ij,k} B_\alpha^i B_\beta^j (p^\sigma B_\sigma^k + \sum_A q_A N_A^k) \\
& + g_{\alpha\sigma} p_{,\beta}^\sigma + G_{ij} B_\alpha^i B_{\sigma,\beta}^j p^\sigma + G_{ij} B_\alpha^i \sum_A (q_{A,\beta} N_A^j + q_A N_{A,\beta}^j) \\
& + g_{\sigma\beta} p_{,\alpha}^\sigma + G_{ij} B_\beta^j B_{\sigma,\alpha}^i p^\sigma + G_{ij} B_\beta^j \sum_A (q_{A,\alpha} N_A^i + q_A N_{A,\alpha}^i) = 0. \quad (1.10)
\end{aligned}$$

Taking covariant derivative of the kind θ with respect to u^μ and using (5), we get

$$\begin{aligned}
z_{|\mu}^i &= p_{|\mu}^\sigma B_\sigma^i + p^\sigma B_{\sigma|\mu}^i + \sum_A (q_{A|\mu} N_A^i + q_A N_{A|\mu}^i) \\
&= p_{|\mu}^\sigma B_\sigma^i + p^\sigma (\Phi_{\sigma\mu}^\pi B_\pi^i + \sum_A \Omega_{A\sigma\mu} N_A^i) + \sum_A q_{A|\mu} N_A^i \\
&+ \sum_A q_A (-e_A g^{\pi\sigma} \Omega_{A\sigma\mu} B_\pi^i + \sum_B \Psi_{BA\mu} N_B^i),
\end{aligned}$$

that is

$$z_{|\mu}^i = P_{\theta\mu}^\pi B_\pi^i + \sum_A Q_{A\mu} N_A^i, \quad (1.11)$$

where

$$P_{\theta\mu}^\pi = p_{|\mu}^\pi + p^\sigma \Phi_{\sigma\mu}^\pi - \sum_A e_A q_A \Omega_{A\sigma\mu} g^{\pi\sigma}, \quad (1.12)$$

$$Q_{A\mu} = p^\sigma \Omega_{A\sigma\mu} + q_{A|\mu} + \sum_B q_B \Psi_{AB\mu}. \quad (1.13)$$

The equation (11) is *derivational formula of the infinitesimal bending field* z^i . So, we have

Theorem 1.1 *If the infinitesimal bending field z^i of the subspace $GR_M \subset GR_N$ is expressed by the tangential and the normal component with respect to the GR_M in the form (9), then the derivation formula (11) is valid, where $|\mu$ is covariant derivative of the kind θ according to u^μ , and P, Q are given in (12) and (13) respectively.*

2 Integrability conditions of derivational formula of the infinitesimal bending field

2.0. Applying to (1.11) covariant derivative of the kind ω with respect to u^ν , we get

$$z_{|\mu|\nu}^i = P_{\theta\mu|\nu}^\pi B_\pi^i + P_{\theta\mu}^\pi B_{|\nu}^i + \sum_A (Q_{A\mu|\nu} N_A^i + Q_{A\mu} N_{A|\nu}^i),$$

and substituting $B_{\pi|\nu}^i$ and $N_{A|\nu}^i$ with respect to (1.5), after arranging one obtains

$$\begin{aligned} z_{\theta|\mu}^i|_{\nu} &= [P_{\theta}^{\pi}|_{\omega}^{\mu} + P_{\theta}^{\sigma}\Phi_{\omega}^{\pi\sigma\nu} - \sum_A e_A Q_{A\mu} g^{\pi\sigma} \Omega_{A\sigma\nu}] B_{\pi}^i \\ &+ \sum_A [P_{\theta}^{\pi}\Omega_{A\pi\nu} + Q_{\theta}^{A\mu}|_{\omega}^{\nu} + \sum_B Q_{\theta}^{B\mu} \Psi_{\omega}^{AB\nu}] N_A^i, \end{aligned} \quad (2.1)$$

where the tensors P, Q are given at (1.12,13). From (1) one gets

$$\begin{aligned} z_{\theta|\mu}^i|_{\nu} - z_{\omega}^i|_{\nu|\mu} &= [P_{\theta}^{\pi}|_{\omega}^{\mu} - P_{\omega}^{\pi}|_{\theta}^{\mu} + P_{\theta}^{\sigma}\Phi_{\omega}^{\pi\sigma\nu} - P_{\omega}^{\sigma}\Phi_{\theta}^{\pi\sigma\mu} \\ &- \sum_A e_A g^{\pi\sigma} (Q_{\theta}^{A\mu} \Omega_{A\sigma\nu} - Q_{\omega}^{A\nu} \Omega_{A\sigma\mu})] B_{\pi}^i \\ &+ \sum_A [P_{\theta}^{\pi}\Omega_{A\pi\nu} - P_{\omega}^{\pi}\Omega_{A\pi\mu} + Q_{\theta}^{A\mu}|_{\omega}^{\nu} - Q_{\omega}^{A\mu}|_{\theta}^{\mu} \\ &+ \sum_B (Q_{\theta}^{B\mu} \Psi_{\omega}^{AB\nu} - Q_{\omega}^{B\nu} \Psi_{\theta}^{AB\mu})] N_A^i. \end{aligned} \quad (2.2)$$

On the other hand applying the Ricci type identities [11,12], we obtain

$$z_{\frac{1}{2}|\mu\nu}^i - z_{\frac{1}{2}|\nu\mu}^i = \frac{R_{1\,pmn}^i}{2} z^p B_{\mu}^m B_{\nu}^n + 2\tilde{\Gamma}_{\mu\nu}^{\pi} z_{\frac{1}{2}}^i|_{\pi}, \quad (2.3a, b)$$

$$z_{\frac{1}{2}|\mu}^i|_{\nu} - z_{\frac{1}{2}|\nu}^i|_{\mu} = R_{3\,p\mu\nu}^i z^p, \quad (2.4)$$

$$z_{\frac{3}{4}|\mu\nu}^i - z_{\frac{3}{4}|\nu\mu}^i = \frac{R_{1\,pmn}^i}{2} z^p B_{\mu}^m B_{\nu}^n \pm 2\tilde{\Gamma}_{\mu\nu}^{\pi} z_{\frac{1}{2}}^i|_{\pi}, \quad (2.5a, b)$$

$$z_{\frac{3}{4}|\mu}^i|_{\nu} - z_{\frac{3}{4}|\nu}^i|_{\mu} = R_{4\,p\mu\nu}^i z^p, \quad (2.6)$$

where [11,12]:

$$R_{1\,jmn}^i = \Gamma_{jm,n}^i - \Gamma_{jn,m}^i + \Gamma_{jm}^p \Gamma_{pn}^i - \Gamma_{jn}^p \Gamma_{pm}^i, \quad (2.7)$$

$$R_{2\,jmn}^i = \Gamma_{mj,n}^i - \Gamma_{nj,m}^i + \Gamma_{mj}^p \Gamma_{np}^i - \Gamma_{nj}^p \Gamma_{mp}^i, \quad (2.8)$$

$$\begin{aligned} R_{3\,j\mu\nu}^i &= (\Gamma_{jm,n}^i - \Gamma_{nj,m}^i + \Gamma_{jm}^p \Gamma_{np}^i - \Gamma_{nj}^p \Gamma_{pm}^i) B_{\mu}^m B_{\nu}^n \\ &+ 2\Gamma_{jm}^i (B_{\mu,\nu}^m - \tilde{\Gamma}_{\nu\mu}^{\pi} B_{\pi}^m), \end{aligned} \quad (2.9)$$

$$\begin{aligned} R_{4\,j\mu\nu}^i &= (\Gamma_{jm,n}^i - \Gamma_{nj,m}^i + \Gamma_{jm}^p \Gamma_{np}^i - \Gamma_{nj}^p \Gamma_{pm}^i) B_{\mu}^m B_{\nu}^n \\ &+ 2\Gamma_{jm}^i (B_{\mu,\nu}^m - \tilde{\Gamma}_{\mu\nu}^{\pi} B_{\pi}^m). \end{aligned} \quad (2.10)$$

The magnitudes $R_{1\,jmn}^i, R_{2\,jmn}^i$ are curvature tensors of the first and the second kind respectively of the space GR_N , while the magnitudes $R_{3\,j\mu\nu}^i, R_{4\,j\mu\nu}^i$ are also

tensors and we called them in [11,12] curvature tensors of the space GR_N with respect to the subspace GR_M .

2.1. The cases (3.a,b) can be written in the form

$$z^i_{|\mu\nu} - z^i_{|\nu\mu} = R^i_{\theta pmn} z^p B^m_\mu B^n_\nu + 2(-1)^\theta \tilde{\Gamma}^\pi_{\mu\nu} z^i_{|\pi}, \theta \in \{1, 2\}. \quad (2.11)$$

Taking in (2) $\omega = \theta \in \{1, 2\}$, we obtain an equation with the same left side as in (11). Substituting $z^i_{|\pi}$ in (11) by virtue of (1.11) and equating the right sides of cited equations, we obtain *the first and the second integrability condition of derivational formula* (1.11) of the infinitesimal bending field z^i of the subspace (for $\theta = 1, \theta = 2$):

$$\begin{aligned} & R^i_{\theta pmn} z^p B^m_\mu B^n_\nu + 2(-1)^\theta \tilde{\Gamma}^\pi_{\mu\nu} (P^\sigma_\theta B^i_\sigma + \sum_A Q_{A\pi} N^i_A) \\ &= [P^\pi_{\theta\mu} |_\nu - P^\pi_{\theta\nu} |_\mu + P^\sigma_{\theta\mu} \Phi^\pi_{\theta\sigma\nu} - P^\sigma_{\theta\nu} \Phi^\pi_{\theta\sigma\mu} \\ &- \sum_A e_A g^{\pi\sigma} (Q_{A\mu} \Omega_{\theta A\sigma\nu} - Q_{A\nu} \Omega_{\theta A\sigma\mu})] B^i_\pi \\ &+ \sum_A [P^\pi_{\theta\mu} \Omega_{\theta A\pi\nu} - P^\pi_{\theta\nu} \Omega_{\theta A\pi\mu} + Q_{\theta A\mu} |_\nu - Q_{\theta A\nu} |_\mu \\ &+ \sum_B (Q_{B\mu} \Psi_{\theta AB\nu} - Q_{B\nu} \Psi_{\theta AB\mu})] N^i_A, \quad \theta = 1, 2. \end{aligned} \quad (2.12)$$

a) Multiplying this equation with $G_{i\lambda} B^i_\lambda$ and using (0.2), (1.1,2), we obtain

$$\begin{aligned} & R_{\theta lpmn} B^l_\lambda z^p B^m_\mu B^n_\nu + 2(-1)^\theta \tilde{\Gamma}^\pi_{\mu\nu} P^\sigma_\pi g_{\lambda\sigma} \\ &= \left[P^\pi_{\theta\mu} |_\nu - P^\pi_{\theta\nu} |_\mu + P^\sigma_{\theta\mu} \Phi^\pi_{\theta\sigma\nu} - P^\sigma_{\theta\nu} \Phi^\pi_{\theta\sigma\mu} - \sum_A e_A g^{\pi\sigma} (Q_{A\mu} \Omega_{\theta A\sigma\nu} - Q_{A\nu} \Omega_{\theta A\sigma\mu}) \right] g_{\lambda\pi}. \end{aligned}$$

Taking into consideration (1.1b) and substituting P, Q according to (1.12,13),

the previous equation becomes

$$\begin{aligned}
 & R_{lpmn} B_\lambda^l z^p B_\mu^m B_\nu^n + 2(-1)^\theta \tilde{\Gamma}_{\mu\nu}^\pi g_{\lambda\sigma} (p_{|\pi}^\sigma + p^\rho \Phi_{\rho\pi}^\sigma - \sum_A e_A q_A \Omega_{A\rho\pi} g^{\sigma\rho}) \\
 &= [p_{|\mu\nu}^\pi + p_{|\nu}^\sigma \Phi_{\sigma\mu}^\pi + p^\sigma \Phi_{\sigma\mu}^\pi |_\nu - \sum_A e_A (q_{A|\nu} \Omega_{A\sigma\mu} + q_A \Omega_{\theta A\sigma\mu} |_\nu) g^{\pi\sigma} \\
 &\quad - p_{|\nu\mu}^\pi - p_{|\mu}^\sigma \Phi_{\theta\sigma\nu}^\pi - p^\sigma \Phi_{\theta\sigma\nu}^\pi |_\mu + \sum_A e_A (q_{A|\mu} \Omega_{A\sigma\nu} + q_A \Omega_{\theta A\sigma\nu} |_\mu) g^{\pi\sigma} \\
 &\quad\quad + (p_{|\mu}^\sigma + p^\rho \Phi_{\rho\mu}^\sigma - \sum_A e_A q_A \Omega_{A\rho\mu} g^{\sigma\rho}) \Phi_{\theta\sigma\nu}^\pi \\
 &\quad\quad - (p_{|\nu}^\sigma + p^\rho \Phi_{\rho\nu}^\sigma - \sum_A e_A q_A \Omega_{A\rho\nu} g^{\sigma\rho}) \Phi_{\theta\sigma\mu}^\pi] g_{\lambda\pi} \\
 &\quad - \sum_A e_A [(p^\sigma \Omega_{A\sigma\mu} + q_{A|\mu} + \sum_B q_B \Psi_{\theta AB\mu}) \Omega_{A\lambda\nu} \\
 &\quad\quad - (p^\sigma \Omega_{A\sigma\nu} + q_{A|\nu} + \sum_B q_B \Psi_{\theta AB\nu}) \Omega_{A\lambda\mu}]. \tag{2.13}
 \end{aligned}$$

Substituting the dummy indices l, p with i, j respectively and z^j according to (1.9), using the Ricci type identity

$$p_{|\mu\nu}^\pi - p_{|\nu\mu}^\pi = \tilde{R}_{\rho\mu\nu}^\pi p^\rho + 2(-1)^\theta \tilde{\Gamma}_{\mu\nu}^\rho p_{|\rho}^\pi, \quad \theta = 1, 2 \tag{2.14}$$

where $\tilde{R}_{\rho\mu\nu}^\pi$ are the corresponding curvature tensors of the subspace (formed by means of $\tilde{\Gamma}$) and denoting

$$\begin{aligned}
 p_\lambda &= g_{\lambda\sigma} p^\sigma, \quad \Phi_{\lambda\rho\pi} = g_{\lambda\sigma} \Phi_{\rho\pi}^\sigma, \\
 \Omega_A^\sigma{}_\mu &= g^{\rho\sigma} \Omega_{A\rho\mu},
 \end{aligned}$$

the equation (13) becomes

$$\begin{aligned}
 & R_{ijmn} B_\lambda^i (p^\sigma B_\sigma^j + \sum_A q_A N_A^j) B_\mu^m B_\nu^n \\
 &+ 2(-1)^\theta \tilde{\Gamma}_{\mu\nu}^\sigma (p^\rho \Phi_{\lambda\rho\sigma} - \sum_A e_A q_A \Omega_{A\lambda\sigma}) \\
 &= p^\sigma (\tilde{R}_{\lambda\sigma\mu\nu} + \Phi_{\theta\lambda\sigma\mu} |_\nu - \Phi_{\theta\lambda\sigma\nu} |_\mu + \Phi_{\theta\sigma\mu}^\rho \Phi_{\lambda\rho\nu} - \Phi_{\theta\sigma\nu}^\rho \Phi_{\lambda\rho\mu}) \\
 &+ \sum_A e_A [q_A (\Phi_{\theta\lambda\sigma\mu} \Omega_A^\sigma{}_\nu - \Phi_{\theta\lambda\sigma\nu} \Omega_A^\sigma{}_\mu - \Omega_{A\lambda\mu} |_\nu + \Omega_{A\lambda\nu} |_\mu) \\
 &\quad\quad + p^\sigma (\Omega_{A\lambda\mu} \Omega_{A\sigma\nu} - \Omega_{A\lambda\nu} \Omega_{A\sigma\mu}) \\
 &\quad\quad + \sum_B q_B (\Omega_{A\lambda\mu} \Psi_{\theta AB\nu} - \Omega_{A\lambda\nu} \Psi_{\theta AB\mu})], \quad \theta = 1, 2. \tag{2.15}
 \end{aligned}$$

b) By multiplying (12) with $G_{i\bar{l}}N_c^l$ and taking into consideration (1.1, 2), one obtains

$$\begin{aligned} & R_{lpmn}N_c^l z^p B_\mu^m B_\nu^n + 2(-1)^\theta \tilde{\Gamma}_{\mu\nu}^\pi Q_{C\pi} e_C \\ &= e_C [P_{\theta\mu}^\pi \Omega_{C\pi\nu} - P_{\theta\nu}^\pi \Omega_{C\pi\mu} + Q_{\theta C\mu|\nu} - Q_{\theta C\nu|\mu} \\ &\quad + \sum_B (Q_{\theta B\mu} \Psi_{CB\nu} - Q_{\theta B\nu} \Psi_{CB\mu})]. \end{aligned}$$

Substituting P_{θ}, Q_{θ} as in the previous case, from here we have

$$\begin{aligned} & R_{ijmn}N_c^i z^j B_\mu^m B_\nu^n \\ &+ 2(-1)^\theta \tilde{\Gamma}_{\mu\nu}^\pi e_C (p^\sigma \Omega_{C\sigma\pi} + q_{C|\pi} + \sum_B q_B \Psi_{CB\pi}) \\ &= e^C \{ (p_{\theta\mu}^\pi + p^\sigma \Phi_{\theta\sigma\mu}^\pi - \sum_A e_A q_A \Omega_{A\sigma\mu} g^{\pi\sigma}) \Omega_{C\pi\nu} \\ &\quad - (p_{\theta\nu}^\pi + p^\sigma \Phi_{\theta\sigma\nu}^\pi - \sum_A e_A q_A \Omega_{A\sigma\nu} g^{\pi\sigma}) \Omega_{C\pi\mu} \\ &\quad + p_{\theta\nu}^\sigma \Omega_{C\sigma\mu} + p_{\theta\mu}^\sigma \Omega_{C\sigma\nu} + q_{C|\mu|\nu} \\ &\quad + \sum_B (q_{B|\nu} \Psi_{CB\mu} + q_B \Psi_{CB\mu|\nu}) \\ &\quad - p_{\theta\mu}^\sigma \Omega_{C\sigma\nu} - p_{\theta\nu}^\sigma \Omega_{C\sigma\mu} - q_{C|\nu|\mu} \\ &\quad - \sum_B (q_{B|\mu} \Psi_{CB\nu} + q_B \Psi_{CB\nu|\mu}) \\ &\quad + \sum_B [(p_{\theta}^\sigma \Omega_{B\sigma\mu} + q_{B|\mu} + \sum_A q_A \Psi_{BA\mu}) \Psi_{CB\nu} \\ &\quad - (p_{\theta}^\sigma \Omega_{B\sigma\nu} + q_{B|\nu} + \sum_A q_A \Psi_{BA\nu}) \Psi_{CB\mu}] \}. \end{aligned}$$

Multiplying the both sides of this equation with $e_C = \pm 1$, and taking into count that

$$\begin{aligned} q_{C|\mu} &= \partial q_C / \partial u^\mu = q_{C,\mu}, \quad \theta = 1, 2, \\ q_{C|\frac{1}{2}\frac{1}{2}\nu} &= (q_{C|\mu})_{,\nu} - \tilde{\Gamma}_{\mu\nu}^\pi q_{C|\pi} = q_{C,\mu\nu} - \tilde{\Gamma}_{\nu\mu}^\pi q_{C,\pi} \end{aligned}$$

from where $q_{C|\mu\nu} - q_{C|\nu\mu} = 2(-1)^\theta \tilde{\Gamma}_{\nu\mu}^\pi q_{C,\pi}$, the previous equation can be written

in the form

$$\begin{aligned}
 & e_C R_{ijmn} N_C^i (p^\sigma B_\sigma^j + \sum_A q_A N_A^j) B_\mu^m B_\nu^n \\
 & + (-1)^\theta \tilde{\Gamma}_{\mu\nu}^\pi (p^\sigma \Omega_{C\sigma\pi} + \sum_B q_B \Psi_{CB\pi}) \\
 = & p^\sigma (\Phi_{\sigma\mu}^\pi \Omega_{C\pi\nu} - \Phi_{\sigma\nu}^\pi \Omega_{C\pi\mu} + \Omega_{C\sigma\mu} |_\nu - \Omega_{C\sigma\nu} |_\mu) \\
 & + \sum_A e_A q_A (\Omega_{C\pi\mu} \Omega_{A\nu}^\pi - \Omega_{C\pi\nu} \Omega_{A\mu}^\pi) \\
 & + \sum_A [p^\sigma (\Omega_{A\sigma\mu} \Psi_{CA\nu} - \Omega_{A\sigma\nu} \Psi_{CA\mu}) \\
 & \quad + q_A (\Psi_{CA\mu} |_\nu - \Psi_{CA\nu} |_\mu) \\
 & + \sum_B q_B (\Psi_{AB\mu} \Psi_{CA\nu} - \Psi_{AB\nu} \Psi_{CA\mu})]. \tag{2.16}
 \end{aligned}$$

2.2 Substituting $\theta = 1$, $\omega = 2$ into (2) and using (4), we obtain the *third integrability condition* of derivational formula (1.11) of z^i :

$$\begin{aligned}
 R_{3p\mu\nu}^i z^p & = [P_{1\mu}^\pi |_\nu - P_{2\nu}^\pi |_\mu + P_{1\mu}^\sigma \Phi_{2\sigma\nu}^\pi - P_{2\nu}^\sigma \Phi_{1\sigma\mu}^\pi \\
 & - \sum_A e_A g^{\pi\sigma} (Q_{A\mu} \Omega_{2A\sigma\nu} - Q_{A\nu} \Omega_{1A\sigma\mu})] B_\pi^i \\
 & + \sum_A [P_{1\mu}^\pi \Omega_{2A\pi\nu} - P_{2\nu}^\pi \Omega_{1A\pi\mu} + Q_{1a\mu} |_\nu - Q_{2A\nu} |_\mu \\
 & + \sum_B (Q_{B\mu} \Psi_{2AB\nu} - Q_{2B\nu} \Psi_{1AB\mu})] N_A^i. \tag{2.17}
 \end{aligned}$$

a) By multiplying the previous equation with $G_{i\lambda} B_\lambda^i$ one obtains

$$\begin{aligned}
 R_{3p\mu\nu}^i B_\lambda^i z^p & = [P_{1\mu}^\pi |_\nu - P_{2\nu}^\pi |_\mu + P_{1\mu}^\sigma \Phi_{2\sigma\nu}^\pi - P_{2\nu}^\sigma \Phi_{1\sigma\mu}^\pi \\
 & - \sum_A e_A g^{\pi\sigma} (Q_{A\mu} \Omega_{2A\sigma\nu} - Q_{2A\nu} \Omega_{1A\sigma\mu})] g_{\lambda\pi}.
 \end{aligned}$$

By substitution of P, Q with respect to (1.12,13), from here it follows that

$$\begin{aligned}
R_{3lp\mu\nu}^i B_\lambda^l z^p &= [p_{1|\mu|2}^\pi + p_{2|\nu|1}^\sigma \Phi_{1\sigma\mu}^\pi + p^\sigma \Phi_{1\sigma\mu|2}^\pi \\
&\quad - \sum_A e_A (q_{A|2} q_{1A\sigma\mu} + q_{A1} \Omega_{A\sigma\mu|2}) g^{\pi\sigma} \\
&\quad \quad - p_{2|\nu|1}^\pi + p_{1|\mu}^\sigma \Phi_{2\sigma\nu}^\pi \\
&\quad + \sum_A e_A (q_{A|1} \Omega_{2A\sigma\nu} + q_{A2} \Omega_{2A\sigma\nu|1}) g^{\pi\sigma} \\
&\quad + (p_{1|\mu}^\sigma + p^\sigma \Phi_{1\rho\mu}^\sigma - \sum_A e_A q_{A1} \Omega_{A\rho\mu} g^{\sigma\rho}) \Phi_{2\sigma\nu}^\pi \\
&\quad - (p_{2|\nu}^\sigma + p^\sigma \Phi_{2\rho\nu}^\sigma - \sum_A e_A q_{A2} \Omega_{A\rho\nu} g^{\sigma\rho}) \Phi_{1\sigma\mu}^\pi] g_{\lambda\pi} \\
&\quad - \sum_A e_A [(p_{1A\sigma\mu}^\sigma \Omega_{A\sigma\mu} + q_{A|1} + \sum_B q_B \Psi_{1AB\mu}) \Omega_{2A\lambda\nu} \\
&\quad \quad - (p_{2A\sigma\nu}^\sigma \Omega_{A\sigma\nu} + q_{A|2} + \sum_B q_B \Psi_{2AB\nu}) \Omega_{1A\lambda\mu}]. \tag{2.18}
\end{aligned}$$

Substituting the dummy indices l, p with i, j respectively and using the Ricci-type identity [11]:

$$p_{1|\mu|2}^\pi - p_{2|\nu|1}^\pi = \tilde{R}_{3\sigma\mu\nu}^\pi p^\sigma, \tag{2.19}$$

where

$$\tilde{R}_{3\beta\mu\nu}^\alpha = \tilde{\Gamma}_{\beta\mu,\nu}^\alpha - \tilde{\Gamma}_{\nu\beta,\mu}^\alpha + \tilde{\Gamma}_{\beta\mu}^\sigma \tilde{\Gamma}_{\nu\sigma}^\alpha - \tilde{\Gamma}_{\nu\beta}^\sigma \tilde{\Gamma}_{\sigma\mu}^\alpha + \tilde{\Gamma}_{\nu\mu}^\sigma (\tilde{\Gamma}_{\sigma\beta}^\alpha - \tilde{\Gamma}_{\beta\sigma}^\alpha) \tag{2.20}$$

is the curvature tensor of the 3rd kind of the subspace, the equation (18) becomes

$$\begin{aligned}
&R_{3ij\mu\nu}^i B_\lambda^j (p^\sigma B_\sigma^i + \sum_A q_A N_A^i) \\
&= p^\sigma (\tilde{R}_{3\lambda\sigma\mu\nu} + \Phi_{1\lambda\sigma\mu|2}^\sigma - \Phi_{2\lambda\sigma\nu|1}^\sigma + \Phi_{1\sigma\mu}^\rho \Phi_{2\lambda\rho\nu}^\sigma - \Phi_{2\sigma\nu}^\rho \Phi_{1\lambda\rho\mu}^\sigma) \\
&\quad + \sum_A e_A [q_A (\Phi_{1\lambda\sigma\mu}^\sigma \Omega_{2A\nu}^\sigma - \Phi_{2\lambda\sigma\nu}^\sigma \Omega_{1A\mu}^\sigma - \Omega_{1A\lambda\mu|2}^\sigma + \Omega_{2A\lambda\nu|1}^\sigma) \\
&\quad \quad + p^\sigma (\Omega_{1A\lambda\mu}^\sigma \Omega_{2A\sigma\nu} - \Omega_{2A\lambda\nu}^\sigma \Omega_{1A\sigma\mu}^\sigma) \\
&\quad \quad + \sum_B q_B (\Omega_{1A\lambda\mu}^\sigma \Psi_{2AB\nu} - \Omega_{2A\lambda\nu}^\sigma \Psi_{1AB\mu})]. \tag{2.21}
\end{aligned}$$

b) Multiplying (17) with $G_{il} N_c^l$, one obtains

$$\begin{aligned}
R_{3lp\mu\nu} N_c^l z^p &= e_c [P_{1\mu}^\pi \Omega_{2C\pi\nu} - P_{2\nu}^\pi \Omega_{1C\pi\mu} + Q_{1c\mu|2} - Q_{2c\nu|1} \\
&\quad + \sum_B (Q_{B\mu} \Psi_{2CB\nu} - Q_{B\nu} \Psi_{1CB\mu})].
\end{aligned}$$

Substituting P, Q using that

$$q_{c|\mu|\nu} - q_{c|\nu|\mu} = 0,$$

and arranging, we get

$$\begin{aligned} & e_C R_{3ij\mu\nu} N_c^i (p^\sigma B_\sigma^j + \sum_A q_A N_A^j) \\ &= p^\sigma (\Phi_{1\sigma\mu}^\pi \Omega_{2c\pi\nu} - \Phi_{2\sigma\nu}^\pi \Omega_{1c\pi\mu} + \Omega_{1c\sigma\mu|\nu} - \Omega_{2c\sigma\nu|\mu}) \\ & \quad + \sum_A e_A q_A (\Omega_{1c\pi\mu} \Omega_{2A\nu}^\pi - \Omega_{2c\pi\nu} \Omega_{1A\mu}^\pi) \\ & + \sum_A [p^\sigma (\Omega_{1A\sigma\mu} \Psi_{2CA\nu} - \Omega_{2A\sigma\nu} \Psi_{1CA\mu}) + q_A (\Psi_{1CA\mu|\nu} - \Psi_{2CA\nu|\mu}) \\ & \quad \sum_B q_B (\Psi_{1AB\mu} \Psi_{2CA\nu} - \Psi_{2AB\nu} \Psi_{1CA\mu})]. \end{aligned} \quad (2.22)$$

2.3. The cases (5a,b) can be given with the equation

$$z_{|\mu\nu}^i - z_{|\nu\mu}^i = R_{\theta-2}^i{}^{pmn} z^p B_\mu^m B_\nu^n + 2(-1)^{\theta-1} \tilde{\Gamma}_{\mu\nu}^\pi z_{|\pi}^i, \quad \theta \in \{3, 4\}. \quad (2.23)$$

Substituting $\theta \in \{3, 4\}$ in (2), we get the equation with the left side as in (21). According to that we get *the 4th and the 5th integrability condition* of the derivation formula (1.11) (for $\theta \in \{3, 4\}$):

$$\begin{aligned} & R_{\theta-2}^i{}^{pmn} z^p B_\mu^m B_\nu^n + 2(-1)^{\theta-1} \tilde{\Gamma}_{\mu\nu}^\pi (P_\pi^\sigma B_\sigma^i + \sum_A Q_{A\pi} N_A^i) \\ &= [P_{|\mu}^\pi|_\nu - P_{|\nu}^\pi|_\mu + P_{|\mu}^\sigma \Phi_{|\mu}^\pi \sigma_\nu - P_{|\nu}^\sigma \Phi_{|\nu}^\pi \sigma_\mu \\ & \quad - \sum_A e_A g^{\pi\sigma} (Q_{A\mu} \Omega_{A\sigma\nu} - Q_{A\nu} \Omega_{A\sigma\mu})] B_\pi^i \\ & + \sum_A [P_{|\mu}^\pi \Omega_{A\pi\nu} - P_{|\nu}^\pi \Omega_{A\pi\mu} + Q_{A\mu|\nu} - Q_{A\nu|\mu}] \\ & + \sum_B (Q_{B\mu} \Psi_{AB\nu} - Q_{B\nu} \Psi_{AB\mu}) N_A^i, \quad \theta \in \{3, 4\}. \end{aligned} \quad (2.24)$$

a) Multiplying this equation with $G_{i\lambda} B_\lambda^i$, we get

$$\begin{aligned} & R_{\theta-2}{}^{lpmn} B_\lambda^l z^p B_\mu^m B_\nu^n + 2(-1)^{\theta-1} \tilde{\Gamma}_{\mu\nu}^\pi P_\pi^\sigma g_{\lambda\sigma} \\ &= [P_{|\mu}^\pi|_\nu - P_{|\nu}^\pi|_\mu + P_{|\mu}^\sigma \Phi_{|\mu}^\pi \sigma_\nu - P_{|\nu}^\sigma \Phi_{|\nu}^\pi \sigma_\mu \\ & \quad - \sum_A e_A g^{\pi\sigma} (Q_{A\mu} \Omega_{A\sigma\nu} - Q_{A\nu} \Omega_{A\sigma\mu})] g_{\lambda\pi}. \end{aligned}$$

from where, as in previous cases,

$$\begin{aligned}
& R_{\theta-2}{}^l{}_{p m n} B_\lambda^l z^p B_\mu^m B_\nu^n + 2(-1)^{\theta-1} \tilde{\Gamma}_{\mu\nu}^\pi g_{\lambda\sigma} (p_\theta^\sigma + p^\rho \Phi_\theta^{\sigma\rho}) \\
& - \sum_A e_A q_A \Omega_{A\rho\pi} g^{\sigma\rho} = [p_\theta^\pi{}_{|\mu\nu} + p_\theta^\sigma{}_{|\nu} \Phi_\theta^\pi{}_{\sigma\mu} + p^\sigma \Phi_\theta^\pi{}_{\sigma\mu}{}_{|\nu} \\
& \quad - \sum_A e_A (q_A{}_{|\nu} \Omega_{A\sigma\mu} + q_A \Omega_{A\sigma\mu}{}_{|\nu}) g^{\pi\sigma} \\
& \quad - p_\theta^\pi{}_{|\nu\mu} - p_\theta^\sigma{}_{|\mu} \Phi_\theta^\pi{}_{\sigma\nu} - p^\sigma \Phi_\theta^\pi{}_{\sigma\nu}{}_{|\mu} \\
& \quad + \sum_A e_A (q_A{}_{|\mu} \Omega_{A\sigma\nu} + q_A \Omega_{A\sigma\nu}{}_{|\mu}) g^{\pi\sigma} \\
& \quad + (p_\theta^\sigma{}_{|\mu} + p^\rho \Phi_\theta^{\sigma\rho} - \sum_A e_A q_A \Omega_{A\rho\mu} g^{\sigma\rho}) \Phi_\theta^\pi{}_{\sigma\nu} \\
& \quad - (p_\theta^\sigma{}_{|\nu} + p^\rho \Phi_\theta^{\sigma\rho} - \sum_A e_A q_A \Omega_{A\rho\nu} g^{\sigma\rho}) \Phi_\theta^\pi{}_{\sigma\mu}] g_{\lambda\pi} \\
& - \sum_A e_A [(p_\theta^\sigma \Omega_{A\sigma\mu} + q_A{}_{|\mu} + \sum_B q_B \Psi_{AB\mu}) \Omega_{A\lambda\nu}] \\
& \quad - (p^\sigma \Omega_{A\sigma\nu} + q_A{}_{|\nu} + \sum_B q_B \Psi_{AB\nu}) \Omega_{A\lambda\mu}].
\end{aligned}$$

According to [12]:

$$p_\theta^\pi{}_{|\mu\nu} - p_\theta^\pi{}_{|\nu\mu} = \tilde{R}_{\theta-2}{}^\pi{}_{\sigma\mu\nu} p^\sigma + 2(-1)^{\theta-1} \tilde{\Gamma}_{\mu\nu}^\sigma p_\theta^\pi{}_{|\sigma}, \theta \in \{3, 4\}, \quad (2.25)$$

the previous equation becomes

$$\begin{aligned}
& R_{\theta-2}{}^i{}_{j m n} B_\lambda^i (p^\sigma B_\sigma^j + \sum_A q_A N_A^j) B_\mu^m B_\nu^n \\
& + 2(-1)^{\theta-1} \tilde{\Gamma}_{\mu\nu}^\pi (p^\sigma \Phi_\theta^{\lambda\sigma\pi} - \sum_A e_A q_A \Omega_{A\lambda\pi}) \\
& = p^\sigma (R_{\theta-2}{}^\lambda{}_{\sigma\mu\nu} + \Phi_\theta^{\lambda\sigma\mu}{}_{|\nu} - \Phi_\theta^{\lambda\sigma\nu}{}_{|\mu} + \Phi_\theta^\rho{}_{\sigma\mu} \Phi_\theta^{\lambda\rho\nu} - \Phi_\theta^\rho{}_{\sigma\nu} \Phi_\theta^{\lambda\rho\mu}) \\
& + \sum_A e_A [q_A (\Phi_\theta^{\lambda\sigma\mu} \Omega_{A\nu}^\sigma - \Phi_\theta^{\lambda\sigma\nu} \Omega_{A\mu}^\sigma \Omega_{A\lambda\mu}{}_{|\nu} - \Omega_{A\lambda\nu}{}_{|\mu}) \\
& \quad + p^\sigma (\Omega_{A\lambda\mu} \Omega_{A\sigma\nu} - \Omega_{A\lambda\nu} \Omega_{A\sigma\mu}) \\
& + \sum_B q_B (\Omega_{A\lambda\mu} \Psi_{AB\nu} - \Omega_{A\lambda\nu} \Psi_{AB\mu})], \quad \theta \in \{3, 4\} \quad (2.26)
\end{aligned}$$

b) Multiplying (23) with $G_{i\bar{l}} N_c^l$, we have

$$\begin{aligned}
& R_{\theta-2}{}^l{}_{p m n} N_c^l z^p B_\mu^m B_\nu^n + 2(-1)^{\theta-1} \tilde{\Gamma}_{\mu\nu}^\pi Q_{c\pi} e_c \\
& = e_c [P_\theta^\pi{}_{\mu} \Omega_{C\pi\nu} - P_\theta^\pi{}_{\nu} \Omega_{C\pi\mu} + Q_{c\mu}{}_{|\nu} - Q_{c\nu}{}_{|\mu}] \\
& + \sum_B (Q_{B\mu} \Psi_{CB\nu} - Q_{B\nu} \Psi_{CB\mu}), \quad \theta \in \{3, 4\}.
\end{aligned}$$

Substituting P_θ, Q_θ , one obtains

$$\begin{aligned}
 & e_c R_{\theta-2}{}^{ijmn} N_c^i z^j B_\mu^m B_\nu^n \\
 & + 2(-1)^{\theta-1} \tilde{\Gamma}_{\mu\nu}^\pi (p^\sigma \Omega_{C\sigma\pi} + q_{C|\pi} + \sum_B q_B \Psi_{CB\pi}) \\
 & = (p_{|\mu}^\pi + p^\sigma \Phi_{\theta\sigma\mu}^\pi - \sum_A e_A q_A \Omega_{A\sigma\mu}^\pi g^{\pi\sigma}) \Omega_{C\pi\nu} \\
 & \quad - (p_{|\nu}^\pi + p^\sigma \Phi_{\theta\sigma\nu}^\pi - \sum_A e_A q_A \Omega_{A\sigma\nu}^\pi g^{\pi\sigma}) \Omega_{C\pi\mu} \\
 & + p_{|\nu}^\sigma \Omega_{C\sigma\mu} + p^\sigma \Omega_{C\sigma\mu|\nu} + q_{C|\mu\nu} + \sum_B (q_{B|\nu} \Psi_{CB\mu} + q_B \Psi_{CB|\nu}) \\
 & - p_{|\mu}^\sigma \Omega_{C\sigma\nu} - p^\sigma \Omega_{C\sigma\nu|\mu} - q_{C|\mu\nu} - \sum_B (q_{B|\nu} \Psi_{CB\mu} + q_B \Psi_{CB|\nu}) \\
 & \quad + \sum_B [(p^\sigma \Omega_{B\sigma\mu} + q_{B|\mu} + \sum_A q_A \Psi_{CB\mu}) \Psi_{CB\nu} \\
 & \quad - (p^\sigma \Omega_{B\sigma\nu} + q_{B|\nu} + \sum_A q_A \Psi_{CB\nu}) \Psi_{CB\mu}].
 \end{aligned}$$

Having in mind that for $\theta \in \{3, 4\}$:

$$q_{C|\mu\nu} - q_{C|\nu\mu} = 2(-1)^{\theta-1} \tilde{\Gamma}_{\mu\nu}^\pi q_{C,\pi}, \quad (2.27)$$

the previous equation, after putting in order, becomes

$$\begin{aligned}
 & e_c R_{\theta-2}{}^{ijmn} N_c^i (P^\sigma B_\sigma^j + \sum_A q_A N_A^j) B_\mu^m B_\nu^n \\
 & + 2(-1)^{\theta-1} \tilde{\Gamma}_{\mu\nu}^\pi (p^\sigma \Omega_{C\sigma\pi} + \sum_B q_B \Psi_{CB\pi}) \\
 & = p^\sigma (\Phi_{\theta\sigma\mu}^\pi \Omega_{C\pi\nu} - \Phi_{\theta\sigma\nu}^\pi \Omega_{C\pi\mu} + \Omega_{C\sigma\mu|\nu} - \Omega_{C\sigma\nu|\mu}) \\
 & \quad + \sum_A e_A q_A (\Omega_{C\pi\mu}^\pi \Omega_{A\nu}^\pi - \Omega_{C\pi\nu}^\pi \Omega_{A\mu}^\pi) \\
 & + \sum_A [p^\sigma (\Omega_{A\sigma\mu} \Psi_{CA\nu} - \Omega_{A\sigma\nu} \Psi_{CA\mu}) + q_A (\Psi_{CA\mu|\nu} - \Psi_{CA\nu|\mu}) \\
 & \quad + \sum_B q_B (\Psi_{AB\mu} \Psi_{CA\nu} - \Psi_{AB\nu} \Psi_{CA\mu})], \quad \theta \in \{3, 4\}.
 \end{aligned} \quad (2.28)$$

2.4. For $\theta = 3$, $\omega = 4$ according to (2) and (6) we get

$$\begin{aligned}
 R_{4^p\mu\nu}^i z^p &= [P_{3^{\mu}}^{\pi}|_{\nu} - P_{4^{\nu}}^{\pi}|_{\mu} + P_{3^{\mu}}^{\sigma}\Phi_{4^{\sigma\nu}}^{\pi} - P_{4^{\nu}}^{\sigma}\Phi_{3^{\sigma\mu}}^{\pi}]B_{\pi}^i \\
 &\quad \sum_A e_A g^{\pi\sigma} (Q_{3^A\mu} \Omega_{4^A\sigma\nu} - Q_{4^A\nu} \Omega_{3^A\sigma\mu}) \\
 &\quad + \sum_A [P_{3^{\mu}}^{\pi}\Omega_{4^A\pi\nu} - P_{4^{\nu}}^{\pi}\Omega_{3^A\pi\mu} + Q_{3^A\mu}|_{\nu} - Q_{4^A\nu}|_{\mu} \\
 &\quad + \sum_B (Q_{3^B\mu} \Psi_{4^A B\nu} - Q_{4^B\nu} \Psi_{3^A B\mu})] N_A^i. \tag{2.29}
 \end{aligned}$$

This is the *6th integrability condition* of the derivational formula (1.11) of the deformation field z^i .

a) Multiplying the previous equation with $G_{\underline{i}\underline{l}} B_{\lambda}^{\underline{l}}$, we get

$$\begin{aligned}
 R_{4^l p\mu\nu} B_{\lambda}^l z^p &= [P_{3^{\mu}}^{\pi}|_{\nu} - P_{4^{\nu}}^{\pi}|_{\mu} + P_{3^{\mu}}^{\sigma}\Phi_{4^{\sigma\nu}}^{\pi} - P_{4^{\nu}}^{\sigma}\Phi_{3^{\sigma\mu}}^{\pi}] \\
 &\quad \sum_A e_A g^{\pi\sigma} (Q_{3^A\mu} \Omega_{4^A\sigma\nu} - Q_{4^A\nu} \Omega_{3^A\sigma\mu}) g_{\lambda\pi} \tag{2.30}
 \end{aligned}$$

From here, analogously to the previous cases, using the Ricci type identity [12]

$$p^{\pi}|_{\mu}|_{\nu} - p^{\pi}|_{\nu}|_{\mu} = \tilde{R}_{4^{\sigma\mu\nu}}^{\pi} p^{\sigma}, \tag{2.31}$$

where

$$R_{4^{\beta\mu\nu}}^{\alpha} = \tilde{\Gamma}_{\beta\mu,\nu}^{\alpha} - \tilde{\Gamma}_{\nu\beta,\mu}^{\alpha} + \tilde{\Gamma}_{\beta\mu}^{\sigma} \tilde{\Gamma}_{\nu\sigma}^{\alpha} - \tilde{\Gamma}_{\nu\beta}^{\sigma} \tilde{\Gamma}_{\sigma\mu}^{\alpha} + \tilde{\Gamma}_{\mu\nu}^{\sigma} (\tilde{\Gamma}_{\sigma\beta}^{\alpha} - \tilde{\Gamma}_{\beta\sigma}^{\alpha}), \tag{2.32}$$

is the 4th kind curvature tensor of a subspace, and from (29) we finally get

$$\begin{aligned}
 &R_{4^i j\mu\nu} B_{\lambda}^i (p^{\sigma} B_{\sigma}^j + \sum_A q_A N_A^j) \\
 &= p^{\sigma} (\tilde{R}_{4^{\lambda\sigma\mu\nu}} + \Phi_{3^{\lambda\sigma\mu}}^{\nu} + \Phi_{4^{\lambda\sigma\nu}}^{\mu} + \Phi_{3^{\sigma\mu}}^{\rho} \Phi_{4^{\lambda\rho\nu}} - \Phi_{4^{\sigma\nu}}^{\rho} \Phi_{3^{\lambda\rho\mu}}) \\
 &\quad + \sum_A e_A [q_A (\Phi_{3^{\lambda\rho\mu}}^{\sigma} \Omega_{4^A\nu}^{\sigma} - \Phi_{4^{\lambda\sigma\nu}}^{\rho} \Omega_{3^A\nu}^{\sigma} - \Omega_{3^A\lambda\sigma\mu}^{\nu} + \Omega_{4^A\lambda\sigma\nu}^{\mu}) \\
 &\quad \quad + p^{\sigma} (\Omega_{3^A\lambda\mu}^{\nu} \Omega_{4^A\sigma\nu} - \Omega_{4^A\lambda\nu}^{\sigma} \Omega_{3^A\sigma\mu}) \\
 &\quad \quad + \sum_B q_B (\Omega_{3^A\lambda\mu}^{\nu} \Psi_{4^A B\nu} - \Omega_{4^A\lambda\nu}^{\sigma} \Psi_{3^A B\mu})]. \tag{2.33}
 \end{aligned}$$

b) Multiplying (29) with $G_{i\bar{l}}N_C^l$ and arranging, we get finally

$$\begin{aligned}
 & e_C R_{ij\mu\nu} N_C^i (p^\sigma B_\sigma^j + \sum_A q_A N_A^j) \\
 = & p^\sigma (\Phi_{3\sigma\mu}^\pi \Omega_{4C\pi\nu} - \Phi_{4\sigma\nu}^\pi \Omega_{3C\pi\mu} + \Omega_{3C\sigma\mu|4\nu} - \Omega_{4C\sigma\nu|3\mu}) \\
 & + \sum_A e_A q_A (\Omega_{3C\pi\mu} \Omega_{4A\nu}^\pi - \Omega_{4C\pi\nu} \Omega_{3A\mu}^\pi) \\
 & \sum_A [p^\sigma (\Omega_{3A\sigma\mu} \Psi_{4CA\nu} - \Omega_{4A\sigma\nu} \Omega_{3CA\mu}) \\
 & \quad + q_A (\Psi_{3CA\mu|4\nu} - \Psi_{4CA\nu|3\mu}) \\
 & + \sum_B q_B (\Psi_{3AB\mu} \Psi_{4CA\nu} - \Psi_{4AB\nu} \Psi_{3CA\mu})] \tag{2.34}
 \end{aligned}$$

From the above exposed, the next theorem is valid:

Theorem 2.1 *If the infinitesimal bending field z^i of the subspace $GR_M \subset GR_N$ is expressed by virtue of tangent and normal component in the form (1.9), then the coefficients p^σ, q_A satisfy the conditions (15), (16), (21), (22), (26), (28), (33), (34).*

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