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On maximum principle for weak subsolutions of degenerate parabolic linear equations

SALVATORE BONAFEDE

Abstract. Sufficient conditions are obtained so that a weak subsolution of (0.1), bounded from above on the parabolic boundary of the cylinder Q , turns out to be bounded from above in Q .

Keywords: maximum principle, weak subsolution, degenerate equation

Classification: 35B50, 35K10, 35K65

1. Introduction

In a recent Note [1] this author has indicated sufficient conditions allowing a weak subsolution ⁽¹⁾ of the parabolic differential equation

$$(0.1) \quad - \sum_1^m i \frac{\partial}{\partial x_i} \left(\sum_1^m j a_{ij} \frac{\partial u}{\partial x_j} + d_i u \right) + \left(\sum_1^m i b_i \frac{\partial u}{\partial x_i} + cu - f \right) + \frac{\partial u}{\partial t} = 0$$

bounded from above on the parabolic boundary of the cylinder Q , to turn out to be bounded from above in Q , assuming the ellipticity condition to be of the following kind:

$$\sum_1^m ij a_{ij}(x, t) \xi_i \xi_j \geq \mu \sum_1^m i \xi_i^2$$

with $\mu = \nu(x)\psi(t)$, ν and ψ satisfactory hypotheses sufficiently general.

Such results do not generally require the subsolutions of (0.1) to have second derivatives with respect to the space variables or the derivative with respect to t .

In this Note similar results are obtained regarding a class of subsolutions less weak as compared to the ones considered in [1], working, however, on more restrictive hypotheses concerning the functions f and ψ .

Moreover, the comparison between these results and the results cited above would require either the functions ψ and ψ^{-1} to be essentially bounded, or, working on more restrictive hypotheses concerning the coefficients of (0.1), $\psi \in C^0([0, T]) \cap L^\infty(0, T)$ and ψ^{-1} to be r -integrable with $r \geq 1$.⁽²⁾ When μ is

⁽¹⁾ see Definition 1, p. 134 of [1].

constant, sufficient conditions for the boundedness of weak subsolutions may be obtained from [2], [6] and [9], whilst the case where μ depends on x and t has been studied by A.V. Ivanov in [4] (see Theorem 5.3) but with a further hypothesis (Condition II) obviously limiting the kind of degeneration (with respect to the variable t) and which has been suppressed in [4] (see the remark at p. 41), assuming, however, that the subsolutions have square-integrable derivative on t .

2. Functional spaces

Let \mathbb{R}^m be the Euclidean m -dimensional space having generical point $x \equiv (x_1, \dots, x_m)$, Ω an open and bounded set of \mathbb{R}^m , T a positive number.

The symbol meas_x (meas) will henceforth indicate the m -dimensional ($m + 1$ -dimensional) LEBESGUE's measure.

If $u(x, t)$ is a function defined in Q and k is a real number, we will indicate with $\Omega(t, k)$, $t \in]0, T[$, the set of points of Ω in which $u(x, t) > k$.

Hypothesis 2.1. Let $\nu(x)$ be a positive function defined in Ω such that:

$$\nu(x) \in L^1(\Omega), \nu^{-1}(x) \in L^1_{\text{loc}}(\Omega).$$

$\tilde{H}^1(\nu, \Omega)$ indicates the completion of $C^1(\bar{\Omega})$ with respect to the norm

$$\|u\|_1 = \left(\int_{\Omega} \left(|u|^2 + \sum_1^m i\nu(x) \left| \frac{\partial u}{\partial x_i} \right|^2 \right) dx \right)^{1/2}.$$

$\tilde{H}^1_0(\nu, \Omega)$ is the closure of $C^\infty_0(\Omega)$ in $\tilde{H}^1(\nu, \Omega)$.

Hypothesis 2.2. Let $\psi(t)$ be a positive function defined in $]0, T[$ such that:

$$\psi(t) \in L^1(0, T).$$

$\tilde{H}^{1,0}(\nu\psi, Q(\tau_1, \tau_2))$ ($0 \leq \tau_1 < \tau_2 \leq T$) stands for the completion of $C^1(\overline{Q(\tau_1, \tau_2)})$ with respect to the norm

$$\|u\|_{1,0,(\tau_1,\tau_2)} = \left(\int_{Q(\tau_1,\tau_2)} \left(|u|^2 + \sum_1^m i\nu(x)\psi(t) \left| \frac{\partial u}{\partial x_i} \right|^2 \right) dx dt \right)^{1/2};$$

$$\|u\|_{1,0} = \|u\|_{1,0,(0,T)}.$$

$\tilde{H}^{1,0}(\nu\psi, Q(\tau_1, \tau_2))$ is a HILBERT space with respect to the norm $\|u\|_{1,0,(\tau_1,\tau_2)}$.

$\tilde{H}^{1,0}_0(\nu\psi, Q(\tau_1, \tau_2))$ ($0 \leq \tau_1 < \tau_2 \leq T$) is the closure of $C^\infty_0(Q(\tau_1, \tau_2))$ in $\tilde{H}^{1,0}(\nu\psi, Q(\tau_1, \tau_2))$.

Finally, we will denote with $\tilde{H}^{1,0,*}(\nu\psi, Q)$ the space of functions $u(x, t)$ belonging to $\tilde{H}^{1,0}(\nu\psi, Q)$, continuous in $[0, T]$ with respect to values in $L^2(\Omega)$.

(2) if $\psi \in C^0(]0, T[) \cap L^1(0, T)$ and $\psi^{-1} \in L^1_{\text{loc}}(0, T)$, the space $C^\infty_0(Q)$ is dense in: $W_\psi = \left\{ w \mid w \in L^2_\psi(0, T; H^1_0(\nu, \Omega)), w_t \in L^2_{1/\psi}(0, T; L^2(\Omega)), w(x, 0) = w(x, T) = 0 \text{ a.e. in } \Omega \right\}$ endowed with the graph norm.

Hypothesis 2.3. *Let us assume that:*

$$\psi, \psi^{-1} \in L^\infty_{\text{loc}}(0, T). \tag{3}$$

Definition 1. We will say that a subsolution of the equation (0.1) is a function $u(x, t) \in \overset{*}{H}^{1,0}(\nu\psi, Q)$ such that

$$(2.1) \quad \int_Q \left\{ \sum_1^m i_j a_{ij} \frac{\partial u}{\partial x_j} \frac{\partial \varphi}{\partial x_i} + \sum_1^m i b_i \frac{\partial u}{\partial x_i} \varphi + cu\varphi + \sum_1^m i d_i u \frac{\partial \varphi}{\partial x_i} - u \frac{\partial \varphi}{\partial t} \right\} dx dt \leq \int_Q f \varphi dx dt$$

for any $\varphi \in C_0^\infty(Q)$ such that $\varphi(x, t) \geq 0$ a.e. in Q .

Definition 2. Given a real number k , if $u \in \tilde{H}^{1,0}(\nu\psi, Q(\tau_1, \tau_2))$ ($0 \leq \tau_1 < \tau_2 \leq T$), we will say that $u(x, t) \leq k$ on $\partial\Omega \times [\tau_1, \tau_2]$ if there exists a sequence $\{u_n\}$ of functions of $C^1(\overline{Q}(\tau_1, \tau_2))$ such that

$$u_n(x, t) \leq k \quad \text{on} \quad \partial\Omega \times [\tau_1, \tau_2]$$

and

$$(2.2) \quad \lim_{n \rightarrow \infty} \|u_n - u\|_{1,0,(\tau_1, \tau_2)} = 0.$$

If k is such that $u(x, t) \leq k$ on $\partial\Omega \times [\tau_1, \tau_2]$, we will say that $u(x, t)$ is bounded from above on $\partial\Omega \times [\tau_1, \tau_2]$. In this case, the symbol $\sup_{[\tau_1, \tau_2]}^* u$ stands for the greatest lower bound of the real numbers k such that $u(x, t) \leq k$ on $\partial\Omega \times [\tau_1, \tau_2]$;

$$\sup^* u = \sup_{[0, T]}^* u.$$

Definition 3. We shall say that a function $u(x, t)$ belonging to $\overset{*}{H}^{1,0}(\nu\psi, Q)$ is bounded from above on $(\Omega \times \{t = 0\}) \cup (\partial\Omega \times [0, T])$ if $u(x, 0)$ is bounded from above in Ω and, also, if $u(x, t)$ is bounded from above on $\partial\Omega \times [0, T]$.

⁽³⁾ hypotheses 2.2 and 2.3 do not imply ψ and ψ^{-1} to be essentially bounded in $]0, T[$; e.g. it will be sufficient to consider:

$$\psi(t) = \begin{cases} \sqrt{t} & 0 < t \leq \frac{1}{2} \\ (\sqrt{1-t})^{-1} & \frac{1}{2} < t \leq 1. \end{cases}$$

3. Hypotheses on coefficients

Let us denote with \mathcal{A} the set of pairs (α^*, α) , with $2 \leq \alpha^*, \alpha \leq +\infty$, such that there exists a positive constant β for which

$$\|u\|_{\alpha^*, \alpha} \leq \beta(\|u\|_{2, \infty} + \|u\|_{1, 0}) \tag{4}$$

for any $u \in L^{2, \infty}(Q) \cap \tilde{H}^{1, 0}(\nu\psi, Q)$. The set \mathcal{A} obviously contains the pair $(2, +\infty)$. Let us indicate with \mathcal{B} the subset of \mathcal{A} formed by the pairs (α^*, α) with $2 < \alpha^*, \alpha < +\infty$.

We will need the following

Hypothesis 3.1. *The set \mathcal{B} is not empty.* (5)

It is therefore reasonable to postulate the following hypotheses on the coefficients of (0.1):

Hypothesis 3.2. *The functions $a_{i,j}, b_i, c, d_i, f$ ($i, j = 1, \dots, m$) are defined and measurable in Q ;*

$$a_{i,j}(\nu\psi)^{-1} \in L^\infty(Q), \quad b_i(\nu\psi)^{-1/2} \in L^{p^*, p}(Q),$$

$$c \in L^{q^*, q}(Q), \quad d_i(\nu\psi)^{-1/2} \in L^{r^*, r}(Q), \quad f \in L^{g^*, g}(Q),$$

where $p^*, p, q^*, q, r^*, r, g^*, g$ are to be such that $2 < 2g^*, 2g < +\infty$

$$\frac{1}{p^*} + \frac{1}{\alpha_1^*} = \frac{1}{2}, \quad \frac{1}{p} + \frac{1}{\alpha_1} = \frac{1}{2}, \quad \frac{1}{q^*} + \frac{2}{\alpha_2^*} = 1$$

$$\frac{1}{q} + \frac{2}{\alpha_2} = 1, \quad \frac{1}{r^*} + \frac{1}{\alpha_3^*} = \frac{1}{2}, \quad \frac{1}{r} + \frac{1}{\alpha_3} = \frac{1}{2}$$

$$\frac{1}{g^*} + \frac{2}{\alpha_4^*} < 1, \quad \frac{1}{g} + \frac{2}{\alpha_4} < 1$$

with $(\alpha_1^*, \alpha_1), (\alpha_2^*, \alpha_2), (\alpha_3^*, \alpha_3)$, belonging to \mathcal{A} and (α_4^*, α_4) belonging to \mathcal{B} .

Moreover, if $p = +\infty$ [$q = +\infty, r = +\infty$] and $p^* < +\infty$ [$q^* < +\infty, r^* < +\infty$], then there exists a function $\eta_1(\sigma)$ [$\eta_2(\sigma), \eta_3(\sigma)$], defined for $\sigma \geq 0$, non decreasing, vanishing for σ approaching zero, having such a property as to give, for almost any t in the interval $]0, T[$:

$$\sum_1^m i \left(\int_E \left(\frac{|b_i(x, t)|}{\sqrt{\nu(x)}} \right)^{p^*} dx \right)^{1/p} \leq \eta_1(\sigma) \sqrt{\psi(t)}$$

$$\left[\left(\int_E (|c(x, t)| - c(x, t))^{q^*} dx \right)^{1/q} \leq \eta_2(\sigma), \right.$$

$$\left. \sum_1^m i \left(\int_E \left(\frac{|d_i(x, t)|}{\sqrt{\nu(x)}} \right)^{r^*} dx \right)^{1/r} \leq \eta_3(\sigma) \sqrt{\psi(t)} \right]$$

(4) if $1 \leq p, q \leq +\infty$, $\|\cdot\|_{p, q, (\tau_1, \tau_2)}$ stands for the norm in $L^{p, q}(Q(\tau_1, \tau_2))$ ($0 \leq \tau_1 < \tau_2 \leq T$); $\|\cdot\|_{p, q} = \|\cdot\|_{p, q, (0, T)}$.

(5) sufficient conditions so that Hypothesis 3.1 holds may be obtained from § 2 of [7].

for all measurable subsets E of Ω such that $\text{meas}_x E \leq \sigma$.

Hypothesis 3.3. *The following inequality results a.e. in Q for all the real numbers $\xi_1, \xi_2, \dots, \xi_m$*

$$\sum_1^m i_j a_{ij}(x, t) \xi_i \xi_j \geq \nu(x) \psi(t) \sum_1^m i \xi_i^2 .$$

Hypothesis 3.4. *There exists a nonnegative constant ϱ :*

$$c - \sum_1^m i \frac{\partial d_i}{\partial x_i} \geq -\varrho$$

in the distributional sense over Q .

In §5 we will prove the following

Theorem. *Let us assume Hypotheses 2.1, 2.2, 2.3, 3.1, 3.2, 3.3, 3.4 hold and let $u(x, t)$ be a subsolution of the equation (0.1) bounded from above on $(\Omega \times \{t = 0\}) \cup (\partial\Omega \times [0, T])$. Then u is bounded from above in Q ; moreover,*

$$\text{ess sup}_Q u \leq e^{\varrho T} \{ \max(0, \text{ess sup}_\Omega u(x, 0), \sup^* u) + \gamma \|f\|_{g^*, g} \} . \tag{6}$$

4. Preliminary lemmas

Lemma 4.1. *Let $u \in \tilde{H}^{1,0}(\nu\psi, Q(\tau_1, \tau_2))$ ($0 \leq \tau_1 < \tau_2 \leq T$) bounded from above on $\partial\Omega \times [\tau_1, \tau_2]$ and $h > \sup^* u$, then there exists a sequence of functions $\{U_\nu\}_\nu$ such that*

$$U_\nu \in C^1(\overline{Q(\tau_1, \tau_2)}), \quad U_\nu(x, t) < h \quad \text{on} \quad \partial\Omega \times [\tau_1, \tau_2]$$

for any $\nu \in \mathbb{N}$ and

$$\lim_{\nu \rightarrow \infty} \|U_\nu - u\|_{1,0,(\tau_1, \tau_2)} = 0 .$$

⁽⁶⁾ having fixed a number

$$l \geq \max(1, T, \beta, \text{meas}_i \Omega, \sum_1^m i \| \frac{b_i(x, t)}{\sqrt{\nu\psi}} \|_{p^*, p}, \|c\|_{q^*, q}, \sum_1^m i \| \frac{d_i(x, t)}{\sqrt{\nu\psi}} \|_{r^*, r})$$

if $p < +\infty$ or $p^* = p = +\infty$, $q < +\infty$ or $q^* = q = +\infty$, $r < +\infty$ or $r^* = r = +\infty$, γ stands for a constant dependent on m, p, q, r, l . If $p = +\infty$ [$q = +\infty, r = +\infty$] and $p^* < +\infty$ [$q^* < +\infty, r^* < +\infty$] γ stands for a constant dependent on $m, p, q, r, l, \eta_1(\sigma)$ [$\eta_2(\sigma), \eta_3(\sigma)$]. The constants dependent on the same arguments will be denoted with the same symbol although their values are different.

It is possible to get, for any $\nu \in \mathbb{N}$, a sequence $\{u_{\nu,n}\}_n$ such that:

$$u_{\nu,n} \in C^1(\overline{Q(\tau_1, \tau_2)}), \quad u_{\nu,n} < h + \frac{1}{\nu} \text{ on } \partial\Omega \times [\tau_1, \tau_2] \text{ for any } n \in \mathbb{N}$$

and

$$\lim_{n \rightarrow \infty} \|u_{\nu,n} - u\|_{1,0,(\tau_1, \tau_2)} = 0.$$

Let us fix $n_\nu \in \mathbb{N}$ such that $\|u_{\nu,n_\nu} - u\|_{1,0,(\tau_1, \tau_2)} < \frac{1}{\nu}$, and assume for any $\nu \in \mathbb{N}$, $U_\nu = u_{\nu,n_\nu} - \frac{1}{\nu}$.

We get $U_\nu \in C^1(\overline{Q(\tau_1, \tau_2)})$, $U_\nu(x, t) < h$ on $\partial\Omega \times [\tau_1, \tau_2]$ for any $\nu \in \mathbb{N}$ and also:

$$\|U_\nu - u\|_{1,0,(\tau_1, \tau_2)} \leq \|u_{\nu,n} - u\|_{1,0,(\tau_1, \tau_2)} + \frac{1}{\nu}(\text{meas } Q)^{1/2}.$$

Lemma 4.2. *Let $u \in \tilde{H}^{1,0}(\nu\psi, Q(\tau_1, \tau_2))$ ($0 \leq \tau_1 < \tau_2 \leq T$) bounded from above on $\partial\Omega \times [\tau_1, \tau_2]$ and $k > \sup_{[\tau_1, \tau_2]} {}^*u$, then $v = u - \min(u, k)$ in $Q(\tau_1, \tau_2)$ belongs to*

$$\tilde{H}^{1,0}(\nu\psi, Q(\tau_1, \tau_2)). \quad (7)$$

Let us fix $k_1 : \sup_{[\tau_1, \tau_2]} {}^*u < k_1 < k$, then (Lemma 4.1) there exists a sequence of functions $\{u_n\}_n$ such that

$$u_n(x, t) \in C^1(\overline{Q(\tau_1, \tau_2)}), \quad u_n(x, t) < k_1 \text{ on } \partial\Omega \times [\tau_1, \tau_2] \text{ for any } n \in \mathbb{N}$$

and (2.2) holds.

As $u_n < k_1$ on $\partial\Omega \times [\tau_1, \tau_2]$, u_n is uniformly continuous in $\overline{Q(\tau_1, \tau_2)}$, it is possible to determine a $\delta = \delta(n)$, $\delta > 0$ such that for any (x, t) belonging to $\overline{Q(\tau_1, \tau_2)}$ with

$$d((x, t), \partial\Omega \times [\tau_1, \tau_2]) < \delta,$$

we get: $u_n(x, t) < k$.

Consequently, assuming $\psi_n = u_n - \min(u_n, k)$ in $\overline{Q(\tau_1, \tau_2)}$, for any $n \in \mathbb{N}$, we get:

$$(4.1) \quad \|\psi_n\|_{1,0,(\tau_1, \tau_2)}^2 \leq 2\|u_n\|_{1,0,(\tau_1, \tau_2)}^2 + 2k^2 \text{meas } Q(\tau_1, \tau_2) \quad \text{and} \\ \text{supp}\{\psi_n(x, t)\} \subset \Omega \times [\tau_1, \tau_2]. \quad (8)$$

We call, for $\mu \in \mathbb{N}$ great enough, $\alpha_\mu(t)$ the characteristic function of the interval $]\tau_1 + \frac{1}{\mu}, \tau_2 - \frac{1}{\mu}[$ and we assume in $\overline{Q(\tau_1, \tau_2)} : \chi_{\mu,n} = \alpha_\mu(t)\psi(x, t)$.

(7) the function v does not generally belong to $\tilde{H}^{1,0}(\nu\psi, Q(\tau_1, \tau_2))$.

(8) if $g : C \rightarrow \mathbb{R}$, we denote with the symbol $\text{supp}\{g\}$ the support of g in C .

The functions $\chi_{\mu,n}, \frac{\partial \chi_{\mu,n}}{\partial x_i}$ ($i = 1, 2, \dots, m$) are bounded and have a compact support in $Q(\tau_1, \tau_2)$ for any $\mu, n \in \mathbb{N}$; moreover,

$$(4.2) \quad \lim_{n \rightarrow \infty} \|\chi_{\mu,n} - \psi_n\|_{1,0,(Q(\tau_1, \tau_2))} = 0.$$

Thus, fixed μ and n , a sequence $\{d_\lambda\}$ of nonnegative equibounded functions of $C_0^\infty(Q(\tau_1, \tau_2))$ converging a.e. in $Q(\tau_1, \tau_2)$ to $\chi_{\mu,n}$ can be constructed via a well-known regularization procedure; ⁽⁹⁾ moreover, also the functions in the sequence $\{\frac{\partial d_\lambda}{\partial x_i}\}_\lambda$ are equibounded in $Q(\tau_1, \tau_2)$ and the sequence converges to $\frac{\partial \chi_{\mu,n}}{\partial x_i}$ a.e. in $Q(\tau_1, \tau_2)$ ($i = 1, 2, \dots, m$).

We deduce from LEBESGUE’s theorem that the function $\chi_{\mu,n}$ belongs to $\overset{0}{\tilde{H}}^{1,0}(\nu\psi, Q(\tau_1, \tau_2))$ for any $\mu, n \in \mathbb{N}$.

Recalling (4.1) and (4.2), it is also proved that, for any $n \in \mathbb{N}$,

$$\psi_n \in \overset{0}{\tilde{H}}^{1,0}(\nu\psi, Q(\tau_1, \tau_2)).$$

From (4.1) we deduce that a subsequence $\{\psi_{n_k}\}_k$ weakly converging in $\overset{0}{\tilde{H}}^{1,0}(\nu\psi, Q(\tau_1, \tau_2))$ can be obtained from the sequence $\{\psi_n\}_n$.

On the other hand, ψ_n converges to v in $L^2(Q(\tau_1, \tau_2))$, so that $v \in \overset{0}{\tilde{H}}^{1,0}(\nu\psi, Q(\tau_1, \tau_2))$. ⁽¹⁰⁾

Remark 4.1. In the particular case where $u \in C^1(\overline{Q(\tau_1, \tau_2)})$ and $k > u$ on $\partial\Omega \times [\tau_1, \tau_2]$, the function $v = u - \min(u, k)$ is the limit in $\overset{0}{\tilde{H}}^{1,0}(\nu\psi, Q(\tau_1, \tau_2))$ of a sequence $\{d_\lambda\}_\lambda$ of nonnegative equibounded functions of $C_0^\infty(Q(\tau_1, \tau_2))$ such that the functions of the sequence $\{\frac{\partial d_\lambda}{\partial x_i}\}_\lambda$ are equibounded ($i = 1, 2, \dots, m$).

Lemma 4.3. *Let us assume that Hypotheses 2.1, 2.2, 2.3, 3.1, 3.2, hold and let $u(x, t)$ be a subsolution of the equation (0.1) bounded from above on $(\Omega \times \{t = 0\}) \cup (\partial\Omega \times [0, T])$. Then, if $0 \leq \tilde{\tau}_1 < \tau < T$ and $k > \sup^* u$, we get*

$$(4.3) \quad \begin{aligned} & \int_{Q(\tilde{\tau}_1, \tau)} \left\{ \sum_1^m ij a_{ij} \frac{\partial v}{\partial x_i} \frac{\partial v}{\partial x_j} + \sum_1^m i b_i \frac{\partial v}{\partial x_i} v + cvv + \right. \\ & \left. + \sum_1^m i d_i u \frac{\partial v}{\partial x_i} \right\} dx dt + \frac{1}{2} \int_\Omega v^2(x, \tau) dx \leq \\ & \leq \int_{Q(\tilde{\tau}_1, \tau)} f v dx dt + \frac{1}{2} \int_\Omega v^2(x, \tilde{\tau}_1) dx \end{aligned}$$

⁽⁹⁾ see, e.g. [3, pp. 109–110].

⁽¹⁰⁾ we get $|v - \psi_{n_k}| \leq |u - u_{n_k}|$ a.e. in $Q(\tau_1, \tau_2)$ for any $k \in \mathbb{N}$.

where $v = u - \min(u, k)$ in Q . (11)

Let $\tilde{\tau}_1, \tau$ be such that $0 < \tilde{\tau}_1 < \tau < T$; setting $\tau_1 = \frac{\tau+T}{2}$, we denote with $C_\tau^\infty(Q)$ the set formed by those nonnegative functions of $C_0^\infty(Q)$ whose support is contained in $Q(0, \tau_1)$. Let $\varphi(x, t)$ be a function of $C_\tau^\infty(Q)$, we extend u, φ and the coefficients of $(0, 1)$ in $\Omega \times]-\infty, +\infty[$, assuming that these functions are equal to zero in those points where they are not defined.

Set $\tau_2 = \frac{T-\tau}{2}$, we then define in $\Omega \times]-\infty, +\infty[$ and for any integer ϱ :

$$\begin{aligned} \Phi_\varrho(x, t) &= \frac{\varrho}{\tau_2} \int_{t-(\tau_2/\varrho)}^t \varphi(x, \lambda) d\lambda, & U_\varrho(x, t) &= \frac{\varrho}{\tau_2} \int_t^{t+(\tau_2/\varrho)} u(x, \lambda) d\lambda, \\ A_{i,\varrho}(x, t) &= \frac{\varrho}{\tau_2} \int_t^{t+(\tau_2/\varrho)} \sum_1^m j a_{ij}(x, \lambda) \frac{\partial u(x, \lambda)}{\partial x_i} d\lambda, \\ B_\varrho(x, t) &= \frac{\varrho}{\tau_2} \int_t^{t+(\tau_2/\varrho)} \sum_1^m i b_i(x, \lambda) \frac{\partial u(x, \lambda)}{\partial x_i} d\lambda, \\ C_\varrho(x, t) &= \frac{\varrho}{\tau_2} \int_t^{t+(\tau_2/\varrho)} c(x, \lambda) u(x, \lambda) d\lambda, & F_\varrho(x, t) &= \frac{\varrho}{\tau_2} \int_t^{t+(\tau_2/\varrho)} f(x, \lambda) d\lambda, \\ D_{i,\varrho}(x, t) &= \frac{\varrho}{\tau_2} \int_t^{t+(\tau_2/\varrho)} d_i(x, \lambda) u(x, \lambda) d\lambda. \end{aligned} \tag{12}$$

From (2.1), in correspondence with $\varphi = \Phi_\varrho(x, t)$, via an exchange in the order of the integrations with respect to t and λ , we get:

$$\begin{aligned} (4.4) \quad \int_Q \left\{ \sum_1^m i A_{i,\varrho} \frac{\partial \varphi}{\partial x_i} + B_\varrho \varphi + C_\varrho \varphi + \sum_1^m i D_{i,\varrho} \frac{\partial \varphi}{\partial x_i} + \frac{\partial U_\varrho}{\partial t} \varphi \right\} dx dt &\leq \\ &\leq \int_Q F_\varrho \varphi dx dt \end{aligned}$$

for all φ belonging to the functional class $C_\tau^\infty(Q)$.

Let $h_2 : \sup^* u < h_2 < k$. Because $h_2 > \sup^* u$, there exists a sequence of functions $\{u_n\}_n$ such that $u_n \in C^1(\overline{Q})$, $u_n < h_2$ on $\partial\Omega \times [0, T]$ and satisfying (2.2). For all pairs of positive integer numbers ϱ and n , we assume:

$$U_{\varrho,n}(x, t) = \frac{\varrho}{\tau_2} \int_t^{t+(\tau_2/\varrho)} u_n(x, \lambda) d\lambda;$$

the function $U_{\varrho,n}(x, t)$ is defined in the closure of the cylinder $Q(0, \tau_1)$ and is therein of class C^1 .

(11) according to Lemma 4.2, $v \in \overset{0}{H}^{1,0}(\nu\psi, Q(\tau_1, \tau_2))$ for any $0 \leq \tau_1 < \tau_2 \leq T$.

(12) we will remark that, because $\nu \in L^1(\Omega)$ and $\psi \in L^1(0, T)$, $a_{ij} \frac{\partial u}{\partial x_i}$ is integrable in $]-\infty, +\infty[$.

Let us now introduce the function $V_{\varrho,n}(x, t)$ defined in Q assuming:

$$V_{\varrho,n}(x, t) = \begin{cases} U_{\varrho,n}(x, t) - \min(U_{\varrho,n}(x, t), k) & \text{in } Q(\tilde{\tau}_1, \tau) \\ 0 & \text{in } Q \setminus Q(\tilde{\tau}_1, \tau). \end{cases}$$

Let $\{\chi_\lambda\}_\lambda$ be the sequence of nonnegative equibounded functions of $C_0^\infty(Q(\tilde{\tau}_1, \tau))$, having partial derivatives with respect to x_i equibounded ($i = 1, 2, \dots, m$), approaching $V_{\varrho,n}$ in $\tilde{H}^{1,0}_0(\nu\psi, Q(\tilde{\tau}_1, \tau))$ (see Remark 4.1). From (4.4), in correspondence with $\varphi = \chi_\lambda$, as λ diverges to $+\infty$, we can deduce the following relation:

$$(4.5) \quad \int_{Q(\tilde{\tau}_1, \tau)} \left\{ \sum_1^m i A_{i,\varrho} \frac{\partial V_{\varrho,n}}{\partial x_i} + B_\varrho V_{\varrho,n} + C_\varrho V_{\varrho,n} + \sum_1^m i D_{i,\varrho} \frac{\partial V_{\varrho,n}}{\partial x_i} + \frac{\partial U_\varrho}{\partial t} V_{\varrho,n} \right\} dx dt \leq \int_{Q(\tilde{\tau}_1, \tau)} F_\varrho V_{\varrho,n} dx dt.$$

Setting, then, in Q :

$$V_\varrho(x, t) = U_\varrho(x, t) - \min(U_\varrho(x, t), k),$$

we get:

$$\|V_\varrho\|_{1,0,(\tilde{\tau}_1, \tau)}^2 \leq \hat{c} \|u\|_{1,0}^2 + k^2 \text{meas } Q, \quad (13) \quad \text{for any } \varrho \in \mathbb{N}.$$

The sequence $\{V_{\varrho,n}\}_n$ converges to V_ϱ in both $\tilde{H}^{1,0}_0(\nu\psi, Q(\tilde{\tau}_1, \tau))$ and $L^{2,\infty}(Q(\tilde{\tau}_1, \tau))$; ⁽¹⁴⁾ accordingly, the function V_ϱ belongs to $\tilde{H}^{1,0}_0(\nu\psi, Q(\tilde{\tau}_1, \tau)) \cap L^{2,\infty}(Q(\tilde{\tau}_1, \tau))$. From (4.5) we deduce, as n goes to $+\infty$, the following:

$$(4.6) \quad \int_{Q(\tilde{\tau}_1, \tau)} \left\{ \sum_1^m i A_{i,\varrho} \frac{\partial V_\varrho}{\partial x_i} + B_\varrho V_\varrho + C_\varrho V_\varrho + \sum_1^m i D_{i,\varrho} \frac{\partial V_\varrho}{\partial x_i} + \frac{\partial U_\varrho}{\partial t} V_\varrho \right\} dx dt \leq \int_{Q(\tilde{\tau}_1, \tau)} F_\varrho V_\varrho dx dt.$$

Let us verify, for example, that:

$$\lim_{n \rightarrow \infty} \int_{Q(\tilde{\tau}_1, \tau)} A_{i,\varrho} \frac{\partial V_{\varrho,n}}{\partial x_i} dx dt = \int_{Q(\tilde{\tau}_1, \tau)} A_{i,\varrho} \frac{\partial V_\varrho}{\partial x_i} dx dt.$$

⁽¹³⁾ the constant \hat{c} depends on $\|\psi\|_{\infty,(\tilde{\tau}_1, \tau)}, \|\psi^{-1}\|_{\infty,(\tilde{\tau}_1, \tau)}$.

⁽¹⁴⁾ we will remark that

$$\|V_{\varrho,n} - V_\varrho\|_{1,0,(\tilde{\tau}_1, \tau)}^2 \leq \hat{c} \|u_n - u\|_{1,0}^2 + o\left(\frac{1}{n}\right) \text{ for any } n \in \mathbb{N}.$$

It will suffice to prove that:

$$\lim_{n \rightarrow \infty} \int_{Q(\tilde{\tau}_1, \tau)} |A_{i, \varrho}| \left| \frac{\partial V_{\varrho, n}}{\partial x_i} - \frac{\partial V_{\varrho}}{\partial x_i} \right| dx dt = 0.$$

We get:

$$\begin{aligned} & \int_{Q(\tilde{\tau}_1, \tau)} |A_{i, \varrho}| \left| \frac{\partial V_{\varrho, n}}{\partial x_i} - \frac{\partial V_{\varrho}}{\partial x_i} \right| dx dt \leq \\ & \leq \hat{c} \left\| \frac{A_{i, \varrho}}{\sqrt{\nu(x)}} \right\|_{2,2,(\tilde{\tau}_1, \tau)} \|V_{\varrho, n} - V_{\varrho}\|_{1,0,(\tilde{\tau}_1, \tau)} \leq \\ & \leq \hat{c} \left\| \frac{a_{i,j}}{\nu\psi} \right\|_{\infty} \cdot \|u\|_{1,0} \cdot \|V_{\varrho, n} - V_{\varrho}\|_{1,0,(\tilde{\tau}_1, \tau)} \quad \text{for all } n \in \mathbb{N}. \end{aligned}$$

We can rewrite (4.6) as follows:

$$\begin{aligned} & \int_{Q(\tilde{\tau}_1, \tau)} \left\{ \sum_1^m i A_{i, \varrho} \frac{\partial V_{\varrho}}{\partial x_i} + B_{\varrho} V_{\varrho} + C_{\varrho} V_{\varrho} + \sum_1^m i D_{i, \varrho} \frac{\partial V_{\varrho}}{\partial x_i} \right\} dx dt + \\ (4.7) \quad & + \frac{1}{2} \int_{\Omega_{\varrho}(\tau, k)} |U_{\varrho}(x, \tau) - k|^2 dx \leq \\ & \leq \int_{Q(\tilde{\tau}_1, \tau)} F_{\varrho} V_{\varrho} dx dt + \frac{1}{2} \int_{\Omega_{\varrho}(\tilde{\tau}_1, k)} |U_{\varrho}(x, \tilde{\tau}_1) - k|^2 dx. \quad (15) \end{aligned}$$

We call $v = u - \min(u, k)$ in Q .

Let us remark now that we get:

$$\|V_{\varrho} - v\|_{1,0,(\tilde{\tau}_1, \tau)}^2 \leq \hat{c} \int_{Q(\tilde{\tau}_1, \tau)} |U_{\varrho} - u|^2 + \nu \sum_1^m i \left| \frac{\partial U_{\varrho, n}}{\partial x_i} - \frac{\partial u}{\partial x_i} \right|^2 dx dt + o\left(\frac{1}{\varrho}\right)$$

for any $\varrho \in \mathbb{N}$; so, V_{ϱ} converges to v in $\tilde{H}^{1,0}(\nu\psi, Q(\tilde{\tau}_1, \tau))$. (16)

Moreover, because U_{ϱ} converges to u in $L^2(\Omega)$ uniformly with respect to $t \in [\tilde{\tau}_1, \tau]$, it is proved that V_{ϱ} converges to v in $L^{2,\infty}(Q(\tilde{\tau}_1, \tau))$.

From (4.7), the conclusion now follows via another passage to the limit.

(15) we will denote with $\Omega_{\varrho}(t, k)$ the set of those points of Ω in which $U_{\varrho}(x, t) > k$.

(16) we get (see [5], p. 85):

$$\lim_{\varrho \rightarrow \infty} \int_{Q(\tilde{\tau}_1, \tau)} |U_{\varrho} - u|^2 + \nu \sum_1^m i \left| \frac{\partial U_{\varrho}}{\partial x_i} - \frac{\partial u}{\partial x_i} \right|^2 dx dt = 0.$$

For example, we prove that:

$$\lim_{\varrho \rightarrow \infty} \int_{Q(\tilde{\tau}_1, \tau)} C_\varrho V_\varrho \, dx \, dt = \int_{Q(\tilde{\tau}_1, \tau)} cuv \, dx \, dt .$$

We get:

$$\begin{aligned} & \left| \int_{Q(\tilde{\tau}_1, \tau)} C_\varrho V_\varrho - cuv \, dx \, dt \right| \leq \\ & \leq \|C_\varrho\|_{\frac{\alpha_2^*}{\alpha_2^*-1}, \frac{\alpha_2}{\alpha_2-1}, (\tilde{\tau}_1, \tau)} \|V_\varrho - v\|_{\alpha_2^*, \alpha_2, (\tilde{\tau}_1, \tau)} + \\ & + \|C_\varrho - cu\|_{\frac{\alpha_2^*}{\alpha_2^*-1}, \frac{\alpha_2}{\alpha_2-1}} \|v\|_{\alpha_2^*, \alpha_2} \leq \\ & \leq \beta \|c\|_{q^*, q} \|u\|_{\alpha_2^*, \alpha_2} (\|V_\varrho - v\|_{1,0, (\tilde{\tau}_1, \tau)} + \|V_\varrho - v\|_{2,\infty, (\tilde{\tau}_1, \tau)}) + \\ & \|C_\varrho - cu\|_{\frac{\alpha_2^*}{\alpha_2^*-1}, \frac{\alpha_2}{\alpha_2-1}} \|v\|_{\alpha_2^*, \alpha_2}, \quad (17) \text{ for all } \varrho \in \mathbb{N} . \end{aligned}$$

If $\tilde{\tau}_1 = 0$, assumed $\tau > 0$, it will suffice to consider $\tau_n = \frac{\tau}{n+1}$ for $n \in \mathbb{N}$. Accordingly, we get:

$$\begin{aligned} & \int_{Q(\tau_n, \tau)} \left\{ \sum_1^m i_j a_{ij} \frac{\partial v}{\partial x_i} \frac{\partial v}{\partial x_j} + \sum_1^m i b_i \frac{\partial v}{\partial x_i} v + cuv + \sum_1^m i d_i u \frac{\partial v}{\partial x_i} \right\} dx \, dt + \\ & + \frac{1}{2} \int_\Omega v^2(x, \tau) \, dx \leq \int_{Q(\tau_n, \tau)} f v \, dx \, dt + \frac{1}{2} \int_\Omega v^2(x, \frac{\tau}{n}) \, dx \text{ for any } n \in \mathbb{N} . \end{aligned}$$

The conclusion will follow via another passage to the limit for n approaching $+\infty$, recalling that the function $v(x, t)$ is continuous in $[0, T]$ to values in $L^2(\Omega)$.

Lemma 4.4. *Let us assume Hypotheses 2.1, 2.2, 2.3, 3.1, 3.2, 3.3, 3.4 with $\varrho = 0$, hold and let $u(x, t)$ be a subsolution of the equation (0.1) bounded from above on $(\Omega \times \{t = 0\}) \cup (\partial\Omega \times [0, T])$.*

Then, if $k > \max(0, \text{ess sup}_\Omega u(x, 0), \sup^ u)$, we get:*

$$\|v\|_{1,0} + \|v\|_{2,\infty} \leq \gamma \|f\psi_k\|_{s^*, s}$$

where $v = u - \min(u, k)$ in Q , ψ_k is the characteristic function of the set of points of Q in which $u(x, t) > k$ and s^*, s are defined by the following formulas:

$$\frac{1}{s^*} + \frac{1}{\alpha_4^*} = 1, \quad \frac{1}{s} + \frac{1}{\alpha_4} = 1 .$$

(17) we will remark that the function which equals $V_\varrho - v$ in $Q(\tilde{\tau}_1, \tau)$ and vanishes in the remaining points belongs to $\overset{0}{\tilde{H}}^{1,0}(\nu\psi, Q) \cap L^{2,\infty}(Q)$ and that C_ϱ converges to cu in $L^{\alpha_2^*/(\alpha_2^*-1)}(Q)$.

According to our hypothesis, we deduce:

$$(4.8) \quad \int_Q c\varphi + \sum_1^m i d_i \frac{\partial \varphi}{\partial x_i} dx dt \geq 0$$

for any $\varphi \in C_0^\infty(Q)$ such that $\varphi(x, t) \geq 0$ a.e. in Q .

From (4.8), via the same procedure adopted in Lemma 4.3, we deduce

$$(4.9) \quad \int_{Q(\tilde{\tau}_1, \tau)} c\sigma + \sum_1^m i d_i \frac{\partial \sigma}{\partial x_i} dx dt \geq 0,$$

for any $\tilde{\tau}_1, \tau : 0 \leq \tilde{\tau}_1 < \tau < T$.

Recalling that $k > \max(0, \text{ess sup}_\Omega u(x, 0), \sup^* u)$, from (4.3) and (4.9) we get:

$$\begin{aligned} & \int_{Q(0, \tau)} \left\{ \sum_1^m i j a_{ij} \frac{\partial v}{\partial x_i} \frac{\partial v}{\partial x_j} + \sum_1^m i b_i \frac{\partial v}{\partial x_i} v + cv^2 + \right. \\ & \left. + \sum_1^m i d_i v \frac{\partial v}{\partial x_i} \right\} dx dt + \frac{1}{2} \int_\Omega v^2(x, \tau) dx \leq \int_{Q(0, \tau)} f v dx dt. \end{aligned}$$

With slight modifications of the procedure followed in Lemma 4.1 of [7] the conclusion easily follows.

5. Proof of Theorem

Let us first examine the particular case where $\varrho = 0$.

There is no loss of generality if we assume that

$$(5.1) \quad \alpha_4^* \left(1 - \frac{1}{g^*}\right) \geq \alpha_4 \left(1 - \frac{1}{g}\right);$$

in fact, the first term of the preceding inequality is greater than 2, so that the inequality will hold by decreasing g .

Let \bar{k} be a number greater than $\max(0, \text{ess sup}_\Omega u(x, 0), \sup^* u)$ and h and k two numbers such that $\bar{k} \leq k < h$.

Assumed that $v = u - \min(u, k)$ in Q , we get:

$$(5.2) \quad \|v\|_{\alpha_4^*, \alpha_4} \geq (h - k) \left(\int_0^\tau [\text{mis}_x \Omega(t, h)]^{\alpha_4 / \alpha_4^*} dt \right)^{1/\alpha_4} = (h - k) \|\psi_h\|_{\alpha_4^*, \alpha_4}.$$

On the other hand (Lemma 4.4 and Hypothesis 3.1)

$$\|v\|_{\alpha_4^*, \alpha_4} \leq \gamma \|f\psi_k\|_{s^*, s},$$

then from (5.2), we get:

$$(5.3) \quad \|\psi_h\|_{\alpha_4^*, \alpha_4} \leq \frac{\gamma}{(h-k)} \|f\|_{g^*, g} \|\psi_k\|_{\lambda^*, \lambda}$$

where

$$\frac{1}{\lambda^*} = \frac{1}{s^*} - \frac{1}{g^*} = 1 - \frac{1}{\alpha_4^*} - \frac{1}{g^*} > \frac{1}{\alpha_4^*},$$

$$\frac{1}{\lambda} = \frac{1}{s} - \frac{1}{g} = 1 - \frac{1}{\alpha_4} - \frac{1}{g} > \frac{1}{\alpha_4}.$$

We get:

$$\frac{\lambda}{\lambda^*} = \frac{1 - \frac{1}{\alpha_4^*} - \frac{1}{g^*}}{1 - \frac{1}{\alpha_4} - \frac{1}{g}} \geq \frac{\frac{\alpha_4}{\alpha_4^*} (1 - \frac{1}{g}) - \frac{1}{\alpha_4^*}}{1 - \frac{1}{\alpha_4} - \frac{1}{g}} = \frac{\alpha_4}{\alpha_4^*},$$

from which, a.e. in $]0, T[$:

$$(5.4) \quad [\text{meas}_x \Omega(t, k)]^{\lambda/\lambda^*} \leq l^{\alpha_4} [\text{meas}_x \Omega(t, k)]^{\alpha_4/\alpha_4^*}.$$

From (5.3) and (5.4) we deduce that

$$\|\psi_h\|_{\alpha_4^*, \alpha_4} \leq \frac{\gamma}{(h-k)} \|f\|_{g^*, g} (\|\psi_k\|_{\alpha_4^*, \alpha_4})^\vartheta,$$

where $\vartheta = \frac{\alpha_4}{2} (1 - \frac{1}{g}) > 1$.

If we assume for any $k \geq \bar{k}$:

$$\eta(k) = \|\psi_k\|_{\alpha_4^*, \alpha_4}$$

then we get (see [8], p. 212):

$$(5.5) \quad \eta(\bar{k} + d) = 0, \text{ where } d = \gamma \|f\|_{g^*, g} \|\psi_{\bar{k}}\|_{\alpha_4^*, \alpha_4}^{\vartheta-1} 2^{\vartheta/(\vartheta-1)}.$$

Remarking that $\|\psi_{\bar{k}}\|_{\alpha_4^*, \alpha_4} \leq l$, from (5.5) we get:

$$u(x, t) \leq \bar{k} + 2^{\vartheta/(\vartheta-1)} l^{\vartheta-1} \gamma \|f\|_{g^*, g}$$

a.e. in Q , from which the proof follows in the case where $\varrho = 0$.

Finally, if ϱ is a nonnegative constant, the proof follows as in Theorem of § 3 of [1].

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