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# SOLVABILITY OF A FORCED AUTONOMOUS DUFFING'S EQUATION WITH PERIODIC BOUNDARY CONDITIONS IN THE PRESENCE OF DAMPING

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Summary. Let  $g \colon \mathbf{R} \to \mathbf{R}$  be a continuous function,  $e \colon [0,1] \to \mathbf{R}$  a function in  $L^2[0,1]$  and let  $c \in \mathbf{R}$ ,  $c \neq 0$  be given. It is proved that Duffing's equation u'' + cu' + g(u) = e(x), 0 < x < 1, u(0) = u(1), u'(0) = u'(1) in the presence of the damping term has at least one solution provided there exists an R > 0 such that  $g(u)u \geqslant 0$  for  $|u| \geqslant R$  and  $\int_0^1 e(x) \, \mathrm{d}x = 0$ . It is further proved that if g is strictly increasing on  $\mathbf{R}$  with  $\lim_{u \to -\infty} g(u) = -\infty$ ,  $\lim_{u \to \infty} g(u) = \infty$  and is Lipschitz continuous with Lipschitz constant  $\alpha < 4\pi^2 + c^2$ , then Duffing's equation given above has exactly one solution for every  $e \in L^2[0,1]$ .

Keywords: Duffing's equation, damping

AMS classification: 34B15, 34C25, 47H15

### 1. Introduction

Let  $g: \mathbf{R} \to \mathbf{R}$  be a continuous function,  $e: [0,1] \to \mathbf{R}$  a function in  $L^2[0,1]$  and let  $c \in \mathbf{R}$ ,  $c \neq 0$  be given. This paper is devoted to the existence of a solution of the forced autonomous Duffing's equation

(1.1) 
$$u'' + cu' + g(u) = e, 0 < x < 1,$$
$$u(0) = u(1), u'(0) = u'(1).$$

We call the equation in (1.1) "autonomous" since the nonlinear function g is independent of x. When g is a function of both the variables x and u, i.e.  $g:[0,1]\times\mathbb{R}\to\mathbb{R}$  is a function satisfying Caratheodory's conditions, the non-autonomous Duffing's

equation

$$u'' + cu' + g(x, u) = e, 0 < x < 1,$$

$$u(0) = u(1), u'(0) = u'(1)$$

has been extensively studied earlier (see e.g. [1], [2], [3], [4], [8], among others). It was shown, for example, by Gupta in [1] that if there exists a  $\varrho > 0$  such that  $g(x, u)u \leq 0$  for a.e.  $x \in [0, 1]$  and all  $u \in \mathbf{R}$  with  $|u| \geq R$  then (1.2) has at least one solution provided  $\int_0^1 e(x) dx = 0$ . In the case when there exists a  $\varrho > 0$  such that  $g(x, u)u \geq 0$  for a.e.  $x \in [0, 1]$  and  $|u| \geq \varrho$ , it was shown in [3] that (1.2) has at least one solution provided  $\int_0^1 e(x) dx = 0$  and  $\lim_{|u| \to \infty} \sup_{\|u\| \to \infty} \frac{g(x, u)}{u}$  is strictly less than

 $4\pi^2 + c^2$ . Now when  $c \neq 0$ , then  $4\pi^2 + c^2 > 4\pi^2$ , which is the second eigenvalue of the linear eigenvalue problem

$$-u'' = \lambda u,$$
(1.3) 
$$u(0) = u(1), \qquad u'(0) = u'(1).$$

It was remarked in [3] that  $\lambda = 0$  is the only eigenvalue of the linear eigenvalue problem when  $c \neq 0$ ,

$$u'' + cu' = \lambda u,$$
(1.4) 
$$u(0) = u(1), \qquad u'(0) = u'(1)$$

to explain that the nonlinearity in g(x,u) can resonate beyond the second eigenvalue  $4\pi^2$  of the linear eigenvalue problem (1.3). Indeed, the author feels that when  $c \neq 0$  and  $g(x,u)u \geqslant 0$  for a.e.  $x \in [0,1]$  and  $|u| \geqslant \varrho$ , then the boundary value problem (1.2) should have at least one solution when  $\int_0^1 e(x) \, \mathrm{d}x = 0$ . But this is not known at this time. The purpose of this paper is to prove this conjecture in the case of the autonomous boundary value problem (1.1) when  $c \neq 0$ . The autonomous problem (1.1) was studied, when  $c \neq 0$ , by Nieto and Rao in [8] in the case when  $g: \mathbf{R} \to \mathbf{R}$  is increasing and  $\lim_{u \to \pm \infty} g(u) = g(\pm \infty)$  exists. But this case was already covered in [1] because then g is bounded on  $\mathbf{R}$  and accordingly,  $\lim_{|u| \to \infty} \frac{g(u)}{u} = 0 < 4\pi^2$ .

Our methods involve using Mawhin's version of the Leray-Schauder continuation theorem and Wirtinger type inequalities to get the needed estimates. We also present some uniqueness results for the boundary value problem (1.1).

#### 2. MAIN RESULTS

Let X, Y denote the Banach spaces X = C[0, 1] and  $Y = L^1[0, 1]$  with their usual norms. Let  $Y_2$  be the subspace of Y spanned by the constant function 1 on [0, 1], i.e.,

$$Y_2 = \{ u \in Y \mid u(x) \equiv c \text{ for a.e. } x \in [0, 1], c \in \mathbb{R} \},$$

and let  $Y_1$  be the subspace of Y such that  $Y = Y_1 \oplus Y_2$ . We note that for  $u \in Y$  we can write

(2.1) 
$$u(x) = \left(u(x) - \int_0^1 u(x) \, \mathrm{d}x\right) + \left(\int_0^1 u(x) \, \mathrm{d}x\right)$$

for  $x \in [0, 1]$ . We define the canonical projection operators  $P: Y \to Y_1, Q: Y \to Y_2$  by

$$P(u)(x) = u(x) - \int_0^1 u(x) dx,$$

$$Q(u) = \int_0^1 u(x) dx$$
(2.2)

for  $u \in Y$ . Clearly, Q = I - P, where I denotes the identity mapping on Y, and the projections P and Q are continuous. Now let  $X_2 = X \cap Y_2$ . Clearly  $X_2$  is a closed subspace of X. Let  $X_1$  be the closed subspace of X such that  $X = X_1 \oplus X_2$ . We note that  $P(X) \subset X_1$ ,  $Q(X) \subset X_2$  and the projections  $P|X: X \to X_1$ ,  $Q|X: X \to X_2$  are continuous. In the following, X, Y, P, Q will refer to the Banach spaces and projections as defined and we will not distinguish between P, P|X (resp. Q, Q|X) and rely on the context for proper meaning.

Also for  $u \in X$ ,  $v \in Y$ , let  $(u, v) = \int_0^1 u(x)v(x) dx$  denote the duality pairing between X and Y. We note that for  $u \in X$ ,  $v \in Y$ , such that u = Pu + Qu, v = Pv + Qv we have

(2.3) 
$$(u, v) = (Pu, Pv) + (Qu, Qv).$$

Let  $c \in \mathbb{R}, c \neq 0$  be given. Define a linear operator  $L: D(L) \subset X \to Y$  by setting

$$(2.4) D(L) = \{ u \in X \mid u'(x) \in AC[0,1], \ u(0) = u(1), \ u'(0) = u'(1) \},$$

and for  $u \in D(L)$ ,

$$(2.5) Lu = u'' + cu'.$$

(Here AC[0,1] denotes the space of real-valued absolutely continuous functions on [0,1].) It is easy to see that L is a linear Fredholm mapping with ker  $L=X_2$ ,  $\operatorname{Im} L=Y_1$ . Further, the mapping  $K:Y_1\to X_1$ , defined for  $u\in Y_1$  by

(2.6) 
$$(Ku)(x) = v(x) - \int_0^1 v(x) \, \mathrm{d}x,$$

where

(2.7) 
$$v(x) = \int_0^x \int_0^\xi e^{c(t-\xi)} u(t) dt d\xi - \frac{e^{-cx} - 1}{c(e^c - 1)} \int_0^1 e^{ct} u(t) dt,$$

(note that we have assumed  $c \neq 0$ ), satisfies the following conditions:

(i) for 
$$u \in Y$$
, we have  $KP(u) \in D(L)$ ,  $LKP(u) = P(u)$ ,

(2.8) (ii) 
$$(KP(u), P(u)) \geqslant -\frac{1}{(4\pi^2 + c^2)} ||P(u)||_{L^2[0,1]}^2$$

Indeed, note that for  $v = KP(u) \in D(L)$ ,

$$(KP(u), P(u)) = (v, Lv) = -\int_0^1 v'^2 \geqslant -\frac{1}{4\pi^2 + c^2} ||Lv||_{L^2[0,1]}^2$$

and so  $(KP(u), P(u)) \ge -\frac{1}{4\pi^2 + c^2} ||P(u)||_{L^2[0,1]}^2$  since

$$\begin{split} \|Lv\|_{L^2[0,1]}^2 &= \int_0^1 (v'' + cv')^2 \, \mathrm{d}x = \int_0^1 \left[ (v'')^2 + 2cv'v'' + c^2(v')^2 \right] \, \mathrm{d}x \\ &= \int_0^1 \left[ (v'')^2 + c^2(v')^2 \right] \, \mathrm{d}x \geqslant (4\pi^2 + c^2) \int_0^1 v'^2 \, \mathrm{d}x. \end{split}$$

Let now  $g: \mathbf{R} \to \mathbf{R}$  be a given continuous function. Let  $N: X \to X \subset Y$  be the non-linear mapping defined by

$$(Nu)(x) = g(u(x)), \qquad x \in [0, 1]$$

for  $u \in X$ . It is then easy to see, using Arzèla-Ascoli theorem, that  $KPN: X \to X_1$  is continuous and compact.

**Theorem 1.** Let  $g: \mathbb{R} \to \mathbb{R}$  be a given continuous function. Let c, a, A, r, R with  $a \leq A, r < 0 < R, c \neq 0$  be such that

$$g(u)\geqslant A \text{ for } u\geqslant R$$
 (2.9) 
$$\text{and}$$
 
$$g(u)\leqslant a \text{ for } u\leqslant r.$$

Then, for every given function  $e(x) \in L^2[0,1]$  with  $a \leq \int_0^1 e(x) dx \leq A$ , Duffing's equation

$$u'' + cu' + g(u) = e, 0 < x < 1,$$

$$(2.10) u(0) = u(1), u'(0) = u'(1)$$

has at least one solution.

Proof. Define functions  $g_1: \mathbf{R} \to \mathbf{R}$  and  $e_1: [0,1] \to \mathbf{R}$  by setting

$$g_1(u) = g(u) - \frac{A+a}{2},$$
  
 $e_1(x) = e(x) - \frac{A+a}{2}.$ 

Then  $g_1: \mathbf{R} \to \mathbf{R}$  is a continuous function and  $e_1: [0,1] \to \mathbf{R}$  is such that  $e_1(x) \in L^2[0,1]$ . Furthermore,

$$g_1(u) \geqslant \frac{1}{2}(A-a) \geqslant 0$$
 for  $u \geqslant R$ ,  
 $g_1(u) \leqslant \frac{1}{2}(a-A) \leqslant 0$  for  $u \leqslant R$ ,

and

$$\frac{1}{2}(a-A)\leqslant \int_0^1 e_1(x)\,\mathrm{d}x\leqslant \frac{1}{2}(A-a).$$

Duffing's equation (2.10) is equivalent to the equation

(2.11) 
$$u'' + cu' + g_1(u) = e_1, \quad 0 < x < 1,$$
$$u(0) = u(1), \quad u'(0) = u'(1).$$

Now, for X = C[0,1] and  $Y = L^1[0,1]$  we consider the Niemytski operator  $N: X \to Y$  defined for  $u \in X$  by

$$(Nu)(x) = g_1(u(x)), \qquad x \in [0,1],$$

and the linear operator  $L: D(L) \subset X \to Y$  defined in (2.4), (2.5).

The equation (2.11) is equivalent to the operator equation

$$(2.12) Lu + Nu = e_1$$

in X. To solve (2.12) it suffices to solve the system of equations

$$Pu + KPNu = KPe_1,$$

$$QNu = Qe_1$$

in X. Indeed, if  $u \in X$  solves (2.13), then  $u \in D(L)$  and

$$LPu + LKPNu = Lu + PNu = LKPe_1 = Pe_1,$$
  
 $QNu = Qe_1,$ 

which gives, on adding, that  $Lu + Nu = e_1$ .

Now, (2.13) is clearly equivalent to the single equation

$$(2.14) Pu + QNu + KPNu = KPe_1 + Qe_1,$$

which has the form of a compact perturbation of the Fredholm operator P of index zero. We can, therefore, apply the version given in [6, Theorem 1, Corollary 1] or [7, Theorem IV.4] or [5] of the Leray-Schauder continuation theorem, which ensures the existence of a solution for (2.14) if the set of all possible solutions of the family of equations

$$(2.15) Pu + (1 - \lambda)Qu + \lambda QNu + \lambda KPNu = \lambda KPe_1 + \lambda Qe_1,$$

 $\lambda \in ]0,1[$ , is a priori bounded independently of  $\lambda$ . Now (2.15) is equivalent to the system of equations

(2.16) 
$$Pu + \lambda KPNu = \lambda KPe_1,$$
$$(1 - \lambda)Qu + \lambda QNu = \lambda Qe_1.$$

Let  $u_{\lambda} \in X$  be a solution of (2.16) for some  $\lambda \in ]0,1[$ , then  $u_{\lambda} \in D(L)$  and

(2.17) 
$$Pu_{\lambda} + \lambda KPNu_{\lambda} = \lambda KPe_{1},$$
$$(1 - \lambda)Qu_{\lambda} + \lambda QNu_{\lambda} = \lambda Qe_{1}.$$

It follows that

$$Lu_{\lambda} + (1 - \lambda)Qu_{\lambda} + Nu_{\lambda} = \lambda e_1,$$

i.e.

$$u_{\lambda}'' + cu_{\lambda}' + (1 - \lambda) \int_{0}^{1} u_{\lambda}(x) dx + g_{1}(u_{\lambda}) = \lambda e_{1},$$

$$(2.18) \qquad u_{\lambda}(0) = u_{\lambda}(1), \quad u_{\lambda}'(0) = u_{\lambda}'(1).$$

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Multiplying the equation in (2.17) by  $u'_{\lambda}$  and integrating over [0, 1] we obtain that

$$c\int_0^1 u_\lambda'^2 = \lambda \int_0^1 e_1(x)u_\lambda'(x)\,\mathrm{d}x,$$

which implies, using the Cauchy-Schwarz inequality, that

$$|c||u_{\lambda}'||_{L^{2}[0,1]} \leqslant ||e_{1}||_{L^{2}[0,1]}.$$

Now, we claim that there exists a  $\xi \in [0, 1]$  such that  $r \leq u_{\lambda}(\xi) \leq R$ . Indeed, suppose that  $u_{\lambda}(x) \geq R$  for every  $x \in [0, 1]$ , then we get from the second equation in (2.17) and our assumptions on  $g_1$  and  $e_1$  that

$$(1-\lambda)R + \lambda \cdot \frac{1}{2}(A-a) \leqslant (1-\lambda)Qu_{\lambda} + \lambda QNu_{\lambda} = \lambda Qe_1 \leqslant \lambda \cdot \frac{1}{2}(A-a),$$

so that  $(1 - \lambda)R \leq 0$ , which is a contradiction since  $\lambda \in ]0, 1[$  and R > 0. Similarly,  $u_{\lambda} \leq r$  for  $x \in [0, 1]$  leads to a contradiction. This proves the claim.

Next it follows that for every  $x \in [0, 1]$ 

$$\begin{split} |u_{\lambda}(x)| &\leqslant \max(-r,R) + \int_{0}^{1} |u'_{\lambda}(x)| \, \mathrm{d}x \\ &\leqslant \max(-r,R) + \|u'_{\lambda}\|_{L^{2}[0,1]} \\ &\leqslant \max(-r,R) + \frac{1}{|c|} \|e_{1}\|_{L^{2}[0,1]} \equiv C. \end{split}$$

Hence

$$||u_{\lambda}||_{X} \leq C.$$

where C is a constant independent of  $\lambda \in ]0,1[$ .

This completes the proof of the theorem.

Corollary 1. Let  $g: \mathbb{R} \to \mathbb{R}$  be a continuous function and let  $c \in \mathbb{R}$ ,  $c \neq 0$  be given. Suppose there exists an R > 0 such that  $g(u)u \geq 0$  for  $u \in \mathbb{R}$ ,  $|u| \geq R$ .

Then for every  $e(x) \in L^2[0,1]$  with  $\int_0^1 e(x) dx = 0$ , Duffing's equation (2.9) has at least one solution.

Proof. The proof follows immediately from Theorem 1 with a = A = 0 and r = -R.

**Theorem 2.** Let  $g: \mathbf{R} \to \mathbf{R}$  be a strictly increasing function with  $\lim_{u \to -\infty} g(u) = -\infty$ ,  $\lim_{u \to \infty} g(u) = \infty$  and let  $c \in \mathbf{R}$ ,  $c \neq 0$ . Suppose that g is a Lipschitz continuous function with a Lipschitz constant  $\alpha$ , i.e.

$$|g(u) - g(v)| \leqslant \alpha |u - v|$$

for  $u, v \in \mathbf{R}$ , with

$$(2.21) \qquad \qquad \alpha < 4\pi^2 + c^2$$

for all  $u \in \mathbf{R}$ .

Then for every  $e \in L^2[0,1]$ , the boundary value problem

(2.22) 
$$u'' + cu' + g(u) = e(x), \quad 0 < x < 1$$
$$u(0) = u(1), \quad u'(0) = u'(1),$$

has exactly one solution u in X = C[0, 1].

**Proof.** Under our assumptions, it is easy to see that there exist a, A, r, R with  $a \le A$ , r < 0 < R such that

$$g(u) \leqslant A \text{ for } u \geqslant R,$$
  
 $g(u) \leqslant a \text{ for } u \leqslant r,$ 

and

$$a \leqslant \int_0^1 e(x) \, \mathrm{d}x \leqslant A.$$

Accordingly, Theorem 1 implies that (2.22) has at least one solution u in X. Let, now,  $u_1, u_2 \in X$  be two different solutions for (2.22). Then

$$(2.23) u_1'' - u_2'' + c(u_1' - u_2') + g(u_1) - g(u_2) = 0, \quad 0 < x < 1.$$

It follows that

$$0 = -\int_0^1 (u_1' - u_2')^2 dx + \int_0^1 (g(u_1) - g(u_2))(u_1 - u_2) dx$$

$$= -\int_0^1 (u_1' - u_2')^2 dx + \int_0^1 |g(u_1) - g(u_2)| |u_1 - u_2| dx$$

$$\geq -\frac{1}{4\pi^2 + c^2} ||Lu_1 - Lu_2||_{L^2[0,1]}^2 + \frac{1}{\alpha} \int_0^1 |g(u_1) - g(u_2)|^2 dx$$

$$= \left(\frac{1}{\alpha} - \frac{1}{4\pi^2 + c^2}\right) \int_0^1 |g(u_1) - g(u_2)|^2 dx,$$

in view of (2.23). Using (2.21), we get that

$$g(u_1(x)) = g(u_2(x))$$

for a.e.  $x \in [0, 1]$ , which implies  $u_1(x) = u_2(x)$  for a.e.  $x \in [0, 1]$ , since g is strictly increasing on  $\mathbb{R}$ . Hence  $u_1(x) = u_2(x)$  for every  $x \in [0, 1]$  since  $u_1$ ,  $u_2$  are continuous in [0, 1].

This completes the proof of the theorem.

Remark 1. Theorem 2 seems to imply that Duffing's equation (2.22) in the presence of the non-zero damping term cu' has a unique solution as long as the non-linearity g(u) does not resonate against too many eigenvalues of the linear eigenvalue problem

$$-u'' = \lambda u, \quad 0 < x < 1,$$
  
 
$$u(0) = u(1), \quad u'(0) = u'(1).$$

Also, it indicates that while the presence of even a small amount of damping gives existence, the presence of large enough damping ensures uniqueness.

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