

Alois Švec

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A PROJECTIVE CHARACTERIZATION OF THE VERONESE SURFACE

ALOIS ŠVEC, BRNO

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There are many metric characterizations of the Veronese surface \mathcal{V} ; see [4], [8] and the literature therein, see also a profound paper [1]. Nevertheless, in the following I try to treat a global purely projective characterization of \mathcal{V} . It seems that there exists no satisfactory local theory of surfaces in $P^4(\mathbb{R})$; the treatise of them in [3] excludes wide classes of surfaces (among them \mathcal{V}). Thus a systematic study of global projective properties of surfaces in $P^n(\mathbb{R})$ is needed; the present paper is an initial first step in this direction.

1. First of all, let us explain what we mean by a Veronese surface $\mathcal{V} \subset P^4(\mathbb{R})$, $P^4(\mathbb{R})$ being the projective 4-dimensional space over reals. Let (x, y, z) be orthonormal coordinates in E^3 and (u_1, \dots, u_5) orthonormal coordinates in E^5 , E^p being the p -dimensional Euclidean space. Denote by $S^s(r) \subset E^{s+1}$ the hypersphere of radius r . The mapping $\sigma: S^2(\sqrt{3}) \rightarrow S^4(1) \subset E^5$ let be given by

$$(1.1) \quad \sigma(x, y, z) = (\alpha yz, \alpha xz, \alpha xy, \frac{1}{2}\alpha(x^2 - y^2), \frac{1}{6}\alpha(x^2 + y^2 - 2z^2))$$

with

$$(1.2) \quad \alpha = \frac{1}{3}\sqrt{3}.$$

To each point $\sigma(p)$, $p \in S^2(\sqrt{3})$ we may associate an orthonormal frame $\{m; v_1, \dots, v_5\}$ in E^5 such that

$$(1.3) \quad \begin{aligned} dm &= \omega^1 v_1 + \omega^2 v_2; \\ dv_1 &= \omega_1^2 v_2 + \alpha \omega^2 v_3 + \alpha \omega^1 v_4 - \omega^1 v_5, \quad dv_2 = -\omega_1^2 v_1 + \alpha \omega^1 v_3 - \alpha \omega^2 v_4 - \omega^2 v_5, \\ dv_3 &= -\alpha \omega^2 v_1 - \alpha \omega^1 v_2 - 2\omega_1^2 v_4, \quad dv_4 = -\alpha \omega^1 v_1 + \alpha \omega^2 v_2 + 2\omega_1^2 v_3, \\ dv_5 &= \omega^1 v_1 + \omega^2 v_2; \end{aligned}$$

for this, see [7] where we have to change the sign of v_5 .

Let $\tau: S^4(1) \rightarrow P^4(\mathbb{R})$ be the usual mapping; then $\tau \circ \sigma: S^2(\sqrt{3}) \rightarrow P^4(\mathbb{R})$ is exactly what we are going to call the *Veronese surface* \mathcal{V} .

2. Let us explain several elementary facts from the theory of Laplace transforms in the projective space; for the hyperbolic case, see [5], Chap. IV.

Given a surface $D \rightarrow P^n(\mathbb{R})$, D being a 2-dimensional manifold, let us suppose that the points of the surface satisfy exactly one hyperbolic partial differential equation of order 2. Then we may choose local coordinates u, v on (a domain U of) D in such a way that our surface is given by $x = x(u, v)$, and we have

$$(2.1) \quad x_{uv} = ax_u + bx_v + cx,$$

the subscripts denoting derivatives. The *Laplace transform* of our surface $x(u, v)$ is a mapping $\tilde{x}: U \rightarrow P^n(\mathbb{R})$, $\tilde{x} = \tilde{x}(u, v)$, such that $\tilde{x}(u, v) \in \{x(u, v), x_u(u, v), x_v(u, v)\}$, $\tilde{x}(u, v) \neq x(u, v)$ for each $(u, v) \in U$, and there is a tangent field $t = t(u, v)$ on U satisfying $t(u, v) \tilde{x}(u, v) \in \{x(u, v), \tilde{x}(u, v)\}$ on U ; by $\{z_1, \dots, z_p\}$, we denote the projective subspace through z_1, \dots, z_p . It is known that our surface $x(u, v)$ has exactly two Laplace transforms

$$(2.2) \quad x_1 = x_v - ax, \quad x_{-1} = x_u - bx.$$

Indeed,

$$(2.3) \quad \frac{\partial}{\partial u} x_1 = bx_1 + hx, \quad \frac{\partial}{\partial v} x_{-1} = ax_{-1} + kx$$

with

$$(2.4) \quad h = c + ab - a_u, \quad k = c + ab - b_v.$$

The functions h, k are the so-called *Laplace-Darboux invariants*. In fact, they are not invariants with respect to the transformation of proportionality factor $x \rightarrow \varrho x'$ and the transformation of parameters $u' = u'(u), v' = v'(v)$, but the so-called *point forms*

$$(2.5) \quad \varphi_1 = h \, du \, dv, \quad \varphi_{-1} = k \, du \, dv$$

are; for their geometrical meaning, see [2].

The Laplace transform $x_1(u, v)$ is a surface if and only if $h \neq 0$ on U ; it satisfies the equation

$$(2.6) \quad x_{1uv} = a_1 x_{1u} + b_1 x_{1v} + c_1 x_1$$

with

$$(2.7) \quad a_1 = a + (\log h)_v, \quad b_1 = b, \quad c_1 = c + h - k - b(\log h)_v; \\ (\log h)_v := h^{-1} h_v.$$

Thus it has once again two Laplace transforms, and they are

$$(2.8) \quad x_2 = (x_1)_1 = x_{1v} - a_1 x_1, \quad (x_1)_{-1} = x_{1u} - b_1 x_1 = hx.$$

For the Laplace transform $x_{-1}(u, v)$, $k \neq 0$ on U means that x_{-1} is a surface. The points x_{-1} satisfy

$$(2.9) \quad x_{-1uv} = a_{-1} x_{-1u} + b_{-1} x_{-1v} + c_{-1} x_{-1}$$

with

$$(2.10) \quad a_{-1} = a, \quad b_{-1} = b + (\log k)_u, \quad c_{-1} = c + k - h - a(\log k)_u;$$

$$(\log k)_u := k^{-1}k_u;$$

and the Laplace transforms of the surface $x_{-1}(u, v)$ are

$$(2.11) \quad x_{-2} = (x_{-1})_{-1} = x_{-1u} - b_{-1}x_{-1}, \quad (x_{-1})_1 = x_{-1v} - a_{-1}x_{-1} = kx.$$

The Laplace-Darboux invariants of (2.6) and (2.9) are

$$(2.12) \quad h_1 = 2h - k - (\log h)_{uv}, \quad k_1 = h; \quad h_{-1} = k,$$

$$k_{-1} = 2k - h - (\log k)_{uv},$$

respectively. Thus we get two new invariant point forms

$$(2.13) \quad \varphi_2 = h_1 du dv, \quad \varphi_{-2} = k_{-1} du dv.$$

We may say that, because of (2.3), φ_1 is associated to the line congruence $\{x, x_1\}$, φ_{-1} to $\{x, x_{-1}\}$, φ_2 to $\{x_1, x_2\}$ and φ_{-2} to $\{x_{-1}, x_{-2}\}$. Of course, $x_2(u, v)$ being a surface, we may construct its Laplace transform $x_3(u, v)$, etc.

Now, let the surface $y: D \rightarrow P^n(\mathbb{R})$ satisfy exactly one elliptic partial differential equation of order 2; we say that it has an *elliptic conjugate net*. It is known that we may choose local coordinates (u, v) in such a way that

$$(2.14) \quad y_{uu} + y_{vv} = Ay_u + By_v + Cy.$$

I did not find this case to be mentioned and studied in the literature, but its theory follows easily. We have to pass to the complexification $P^n(\mathbb{C})$ and complexify the tangent bundle of D . Using the complex coordinate $z = u + iv$ and the usual vector fields

$$(2.15) \quad \frac{\partial}{\partial z} = \frac{1}{2} \left(\frac{\partial}{\partial u} - i \frac{\partial}{\partial v} \right), \quad \frac{\partial}{\partial \bar{z}} = \frac{1}{2} \left(\frac{\partial}{\partial u} + i \frac{\partial}{\partial v} \right)$$

on U , (2.14) may be rewritten as

$$(2.16) \quad y_{z\bar{z}} = \mathcal{A}y_z + \bar{\mathcal{A}}y_{\bar{z}} + \mathcal{C}y; \quad \mathcal{A} = \frac{1}{4}(A + iB), \quad \mathcal{C} = \frac{1}{4}C.$$

The Laplace transforms are then

$$(2.17) \quad y_1 = y_z - \mathcal{A}y, \quad y_{-1} = y_{\bar{z}} - \bar{\mathcal{A}}y$$

with

$$(2.18) \quad \frac{\partial}{\partial z} y_1 = \bar{\mathcal{A}}y_1 + Hy, \quad \frac{\partial}{\partial \bar{z}} y_{-1} = \mathcal{A}y_{-1} + Ky;$$

$$(2.19) \quad H = \mathcal{C} + \mathcal{A}\bar{\mathcal{A}} - \mathcal{A}_z, \quad K = \mathcal{C} + \mathcal{A}\bar{\mathcal{A}} - \bar{\mathcal{A}}_{\bar{z}} = \bar{H}.$$

From (2.15), we see that

$$(2.20) \quad y_{-1} = \bar{y}_1.$$

Further, our point forms are

$$(2.21) \quad \psi_1 = H dz d\bar{z}, \quad \psi_{-1} = K dz d\bar{z} = \bar{\psi}_1.$$

Let us suppose $H \neq 0$ on $U \subset D$; then $y_1(u, v)$ is a surface and $y_{-1}(u, v)$ is a surface as well. Then we get the second Laplace transforms

$$(2.22) \quad y_2 = y_{1z} - \mathcal{A}_1 y_1, \quad y_{-2} = y_{-1z} - \bar{\mathcal{A}}_1 y_{-1}; \quad \mathcal{A}_1 = \mathcal{A} + (\log H)_{\bar{z}};$$

see (2.7)–(2.11). Thus again, using (2.20),

$$(2.23) \quad y_{-2} = \bar{y}_2,$$

and we get the point forms

$$(2.24) \quad \psi_2 = H_1 dz d\bar{z}, \quad \psi_{-2} = K_{-1} dz d\bar{z};$$

$$H_1 = 2H - K - (\log H)_{z\bar{z}}, \quad K_{-1} = 2K - H - (\log K)_{z\bar{z}} = \bar{H}_1$$

satisfying, because of (2.19),

$$(2.25) \quad \psi_{-2} = \bar{\psi}_2.$$

3. Consider a surface $D \rightarrow P^4(\mathbb{R})$, D being a 2-dimensional manifold; we are going to restrict ourselves to its coordinate neighborhood U . Let us suppose that our surface carries exactly one elliptic conjugate net and its first and second Laplace transforms exist.

It follows easily that our surface is not contained in a $P^3(\mathbb{R})$. To each point m of our surface (in U), let us associate a frame $\{m_0, \dots, m_4\}$ consisting of analytic points m_i such that the geometric point m_0 coincides with m and m_1, m_2 are situated in the tangent plane of our surface at m . Further, let

$$(3.1) \quad \det \|m_0, \dots, m_4\| = 1.$$

We have the fundamental equations

$$(3.2) \quad dm_0 = \omega_0^0 m_0 + \omega^1 m_1 + \omega^2 m_2, \quad dm_i = \omega_i^0 m_0 + \dots + \omega_i^4 m_4$$

$$(i = 1, \dots, 4)$$

with the usual integrability conditions

$$(3.3) \quad d\omega_i^j = \omega_i^k \wedge \omega_k^j \quad (i, j = 0, \dots, 4);$$

of course,

$$(3.4) \quad \omega^1 := \omega_0^1, \quad \omega^2 := \omega_0^2; \quad \omega_0^3 = \omega_0^4 = 0.$$

From (3.1), we get

$$(3.5) \quad \omega_0^0 + \dots + \omega_4^4 = 0.$$

Let us choose the frames in such a way that

$$(3.6) \quad M_1 := m_1 + im_2, \quad M_2 := m_4 + im_3$$

are the first and the second Laplace transforms of our surface, respectively. The equation (3.2₁) may be written as

$$(3.7) \quad dm_0 = \omega_0^0 m_0 + \tau^1 M_1 + \bar{\tau}^1 \bar{M}_1$$

with

$$(3.8) \quad \tau^1 = \frac{1}{2}(\omega^1 - i\omega^2).$$

Then it is easy to see that

$$(3.9) \quad \begin{aligned} dM_1 &= (\omega_1^0 + i\omega_2^0) m_0 + \frac{1}{2}\{\omega_1^1 + \omega_2^2 + i(\omega_2^1 - \omega_1^2)\} M_1 + \\ &\quad + \frac{1}{2}\{\omega_1^1 - \omega_2^2 + i(\omega_2^1 + \omega_1^2)\} \bar{M}_1 + \frac{1}{2}\{\omega_1^4 + \omega_2^3 + i(\omega_2^4 - \omega_1^3)\} M_2 + \\ &\quad + \frac{1}{2}\{\omega_1^4 - \omega_2^3 + i(\omega_2^4 + \omega_1^3)\} \bar{M}_2, \\ dM_2 &= (\omega_4^0 + i\omega_3^0) m_0 + \frac{1}{2}\{\omega_4^1 + \omega_3^2 + i(\omega_3^1 - \omega_4^2)\} M_1 + \\ &\quad + \frac{1}{2}\{\omega_4^1 - \omega_3^2 + i(\omega_3^1 + \omega_4^2)\} \bar{M}_1 + \frac{1}{2}\{\omega_4^4 + \omega_3^3 + i(\omega_3^4 - \omega_4^3)\} M_2 + \\ &\quad + \frac{1}{2}\{\omega_4^4 - \omega_3^3 + i(\omega_3^4 + \omega_4^3)\} \bar{M}_2. \end{aligned}$$

The definition of M_1 and M_2 yields

$$(3.10) \quad \omega_1^1 - \omega_2^2 + i(\omega_2^1 + \omega_1^2) = 0, \quad \omega_1^4 - \omega_2^3 + i(\omega_2^4 + \omega_1^3) = 0,$$

$$(3.11) \quad \begin{aligned} (\omega_1^0 + i\omega_2^0) \wedge \tau^1 &= \{\omega_1^4 + \omega_2^3 + i(\omega_2^4 - \omega_1^3)\} \wedge \tau^1 = 0, \\ (\omega_4^0 + i\omega_3^0) \wedge \tau^1 &= \{\omega_4^1 - \omega_3^2 + i(\omega_3^1 + \omega_4^2)\} \wedge \tau^1 = \\ &= \{\omega_4^4 - \omega_3^3 + i(\omega_3^4 + \omega_4^3)\} \wedge \tau^1 = 0. \end{aligned}$$

Thus

$$(3.12) \quad \omega_1^1 - \omega_2^2 = \omega_2^1 + \omega_1^2 = \omega_1^4 - \omega_2^3 = \omega_2^4 + \omega_1^3 = 0,$$

and there are real-valued functions A_1, \dots, F_2 such that

$$(3.13) \quad \omega_1^4 + \omega_2^3 = 2(A_1\omega^1 + A_2\omega^2), \quad \omega_2^4 - \omega_1^3 = 2(A_2\omega^1 - A_1\omega^2),$$

$$\omega_4^1 - \omega_3^2 = 2(B_1\omega^1 + B_2\omega^2), \quad \omega_3^1 + \omega_4^2 = 2(B_2\omega^1 - B_1\omega^2),$$

$$(3.14) \quad \omega_4^4 - \omega_3^3 = C_1\omega^1 + C_2\omega^2, \quad \omega_3^4 + \omega_4^3 = C_2\omega^1 - C_1\omega^2,$$

$$(3.15) \quad \omega_4^0 = E_1\omega^1 + E_2\omega^2, \quad \omega_3^0 = E_2\omega^1 - E_1\omega^2,$$

$$\omega_1^0 = F_1\omega^1 - F_2\omega^2, \quad \omega_2^0 = F_2\omega^1 + F_1\omega^2.$$

Let the functions G_1, \dots, H_2 be defined by

$$(3.16) \quad \omega_4^1 + \omega_3^2 = (G_1 + H_1)\omega^1 + (G_2 - H_2)\omega^2,$$

$$\omega_3^1 - \omega_4^2 = (G_2 + H_2)\omega^1 + (H_1 - G_1)\omega^2.$$

Then the system (3.12)–(3.14), (3.16) is equivalent to the system

$$(3.17) \quad \begin{aligned} \omega_2^2 &= \omega_1^1, \quad \omega_2^1 = -\omega_1^2, \\ \omega_4^4 &= \omega_3^3 + C_1\omega^1 + C_2\omega^2, \quad \omega_4^3 = -\omega_3^4 + C_2\omega^1 - C_1\omega^2, \\ \omega_1^4 &= A_1\omega^1 + A_2\omega^2, \quad \omega_2^3 = A_1\omega^1 + A_2\omega^2, \end{aligned}$$

$$\begin{aligned}
\omega_1^3 &= -A_2\omega^1 + A_1\omega^2, & \omega_2^4 &= A_2\omega^1 - A_1\omega^2, \\
\omega_4^1 &= (B_1 + G_1 + H_1)\omega^1 + (B_2 + G_2 - H_2)\omega^2, \\
\omega_3^2 &= (G_1 + H_1 - B_1)\omega^1 + (G_2 - H_2 - B_2)\omega^2, \\
\omega_3^1 &= (G_2 + H_2 + B_2)\omega^1 + (H_1 - G_1 - B_1)\omega^2, \\
\omega_4^2 &= (B_2 - G_2 - H_2)\omega^1 + (G_1 - H_1 - B_1)\omega^2.
\end{aligned}$$

Thus our starting point are the equations (3.17) + (3.15). Let us define

$$\begin{aligned}
(3.18) \quad \mathcal{D}F_1 &= dF_1 + F_1(\omega_0^0 - \omega_1^1), & \mathcal{D}F_2 &= dF_2 + F_2(\omega_0^0 - \omega_1^1), \\
\mathcal{D}A_1 &= dA_1 + A_1(\omega_0^0 - 2\omega_1^1 + \omega_3^3) - A_2(2\omega_1^2 + \omega_3^4), \\
\mathcal{D}A_2 &= dA_2 + A_2(\omega_0^0 - 2\omega_1^1 + \omega_3^3) + A_1(2\omega_1^2 + \omega_3^4), \\
\mathcal{D}E_1 &= dE_1 + E_1(2\omega_0^0 - \omega_1^1 - \omega_3^3) + E_2(\omega_3^4 - \omega_1^2), \\
\mathcal{D}E_2 &= dE_2 + E_2(2\omega_0^0 - \omega_1^1 - \omega_3^3) - E_1(\omega_3^4 - \omega_1^2), \\
\mathcal{D}B_1 &= dB_1 + B_1(\omega_0^0 - \omega_3^3) - B_2(2\omega_1^2 - \omega_3^4), \\
\mathcal{D}B_2 &= dB_2 + B_2(\omega_0^0 - \omega_3^3) + B_1(2\omega_1^2 - \omega_3^4), \\
\mathcal{D}G_1 &= dG_1 + G_1(\omega_0^0 - \omega_3^3) + G_2\omega_3^4, & \mathcal{D}G_2 &= dG_2 + G_2(\omega_0^0 - \omega_3^3) - G_1\omega_3^4, \\
\mathcal{D}H_1 &= dH_1 + H_1(\omega_0^0 - \omega_3^3) + H_2(2\omega_1^2 + \omega_3^4), \\
\mathcal{D}H_2 &= dH_2 + H_2(\omega_0^0 - \omega_3^3) - H_1(2\omega_1^2 + \omega_3^4), \\
\mathcal{D}C_1 &= dC_1 + C_1(\omega_0^0 - \omega_1^1) - C_2(\omega_1^2 - 2\omega_3^4), \\
\mathcal{D}C_2 &= dC_2 + C_2(\omega_0^0 - \omega_1^1) + C_1(\omega_1^2 - 2\omega_3^4).
\end{aligned}$$

Then the differential consequences of (3.17) + (3.15) are

$$(3.19) \quad \mathcal{D}F_1 \wedge \omega^1 - \mathcal{D}F_2 \wedge \omega^2 = 0, \quad \mathcal{D}F_2 \wedge \omega^1 + \mathcal{D}F_1 \wedge \omega^2 = 0,$$

$$(3.20) \quad \begin{aligned} \mathcal{D}A_1 \wedge \omega^1 + \mathcal{D}A_2 \wedge \omega^2 &= (A_1C_2 - A_2C_1)\omega^1 \wedge \omega^2, \\ \mathcal{D}A_2 \wedge \omega^1 - \mathcal{D}A_1 \wedge \omega^2 &= (A_1C_1 + A_2C_2)\omega^1 \wedge \omega^2, \end{aligned}$$

$$(3.21) \quad \begin{aligned} (\mathcal{D}G_1 + \mathcal{D}H_1) \wedge \omega^1 + (\mathcal{D}G_2 - \mathcal{D}H_2) \wedge \omega^2 &= \\ = \{C_1(B_2 + G_2) - C_2(B_1 + G_1)\} \omega^1 \wedge \omega^2, \\ (\mathcal{D}G_2 + \mathcal{D}H_2) \wedge \omega^1 - (\mathcal{D}G_1 - \mathcal{D}H_1) \wedge \omega^2 &= \\ = \{C_1(B_1 - G_1) + C_2(B_2 - G_2)\} \omega^1 \wedge \omega^2, \end{aligned}$$

$$(3.22) \quad \begin{aligned} \mathcal{D}E_1 \wedge \omega^1 + \mathcal{D}E_2 \wedge \omega^2 &= 2(E_2C_1 - E_1C_2 - F_1G_2 - F_2G_1)\omega^1 \wedge \omega^2, \\ \mathcal{D}E_2 \wedge \omega^1 - \mathcal{D}E_1 \wedge \omega^2 &= 2(F_1G_1 - F_2G_2)\omega^1 \wedge \omega^2, \\ \mathcal{D}B_1 \wedge \omega^1 + \mathcal{D}B_2 \wedge \omega^2 &= \{C_1(B_2 + G_2) - C_2(B_1 + G_1) - E_2\} \omega^1 \wedge \omega^2, \\ \mathcal{D}B_2 \wedge \omega^1 - \mathcal{D}B_1 \wedge \omega^2 &= \{C_2(G_1 - B_1) + C_2(G_2 - B_2) + E_1\} \omega^1 \wedge \omega^2, \\ \mathcal{D}C_1 \wedge \omega^1 + \mathcal{D}C_2 \wedge \omega^2 &= 4(A_2B_1 - A_1B_2)\omega^1 \wedge \omega^2, \\ \mathcal{D}C_2 \wedge \omega^1 - \mathcal{D}C_1 \wedge \omega^2 &= \{4(A_1B_1 + A_2B_2) - C_1^2 - C_2^2\} \omega^1 \wedge \omega^2. \end{aligned}$$

From (3.9) and (3.17) we obtain

$$(3.23) \quad \begin{aligned} dM_1 &= F\bar{\tau}^1 m_0 + \tau_1^1 M_1 + A\tau^1 M_2, \\ dM_2 &= E\tau^1 m_0 + (G\tau^1 + H\bar{\tau}^1) M_1 + B\tau^1 \bar{M}_1 + \tau_3^3 M_2 + C\tau^1 M_2 \end{aligned}$$

with

$$(3.24) \quad A = 2(A_1 + iA_2), \dots, H = 2(H_1 + iH_2),$$

$$(3.25) \quad \tau_1^1 = \omega_1^1 - i\omega_1^2, \quad \tau_3^3 = \omega_3^3 + i\omega_3^4 + (C_1 - iC_2)\bar{\tau}^1.$$

The geometrical points m_0, M_1, M_2 are fixed; nevertheless, we may change their factors of proportionality, i.e., choose other analytic points n_0, N_1, N_2 by

$$(3.26) \quad m_0 = Rn_0, \quad M_1 = SN_1, \quad M_2 = TN_2$$

with R an \mathbb{R} -valued and S, T a \mathbb{C} -valued function, respectively. Further, because of (3.1),

$$(3.27) \quad RS\bar{S}T\bar{T} = 1.$$

Then the equations (3.7) + (3.23) become

$$(3.28) \quad \begin{aligned} dn_0 &= \varphi_0^0 n_0 + \varrho^1 N_1 + \bar{\varrho}^1 \bar{N}_1, \\ dN_1 &= F^* \bar{\varrho}^1 n_0 + \varrho_1^1 N_1 + A^* \varrho^1 N_2, \\ dN_2 &= E^* \varrho^1 n_0 + (G^* \varrho^1 + H^* \bar{\varrho}^1) N_1 + B^* \varrho^1 \bar{N}_1 + \varrho_3^3 N_2 + C^* \varrho^1 \bar{N}_2, \end{aligned}$$

and we have

$$(3.29) \quad \begin{aligned} \varrho^1 &= R^{-1} S \tau^1, \quad F^* = R^2 (S\bar{S})^{-1} F, \quad A^* = RS^{-2} TA, \\ E^* &= R^2 S^{-1} T^{-1} E, \quad G^* = RT^{-1} G, \quad H^* = RS\bar{S}^{-1} T^{-1} H, \\ B^* &= RS^{-1} \bar{S} T^{-1} B, \quad C^* = RS^{-1} T^{-1} \bar{C}. \end{aligned}$$

Consequently,

$$(3.30) \quad F^* \varrho^1 \bar{\varrho}^1 = F \tau^1 \bar{\tau}^1, \quad A^* H^* \varrho^1 \bar{\varrho}^1 = AH \tau^1 \bar{\tau}^1,$$

and we see immediately that we get the invariant point forms

$$(3.31) \quad \psi_1 = F \tau^1 \bar{\tau}^1, \quad \psi_2 = AH \tau^1 \bar{\tau}^1;$$

the point forms ψ_{-1} and ψ_{-2} are given by (2.21₃) and (2.25), respectively.

Theorem. Let $D \subset \mathbb{R}^2$ be a bounded domain, ∂D its boundary. Let $m: D \rightarrow P^4(\mathbb{R})$ be a surface, and let us suppose (i) $m(D)$ has exactly one elliptic conjugate net and its Laplace transforms $M_1, M_2: D \rightarrow P^4(\mathbb{C})$ exist; (ii) for the point forms ψ_1 and $\psi_{-1} = \bar{\psi}_1$, we have

$$(3.32) \quad \psi_1 = \psi_{-1};$$

(iii) the (now real) point form ψ_1 is negative definite, ψ_2 does not vanish, and the Gauss curvature κ of $|\psi_1|$ satisfies

$$(3.33) \quad \kappa > \frac{12}{5} \left(1 - \frac{2}{3} \sqrt{3}\right) \doteq -0.371 \quad \text{on } D, \quad d\kappa = 0 \quad \text{on } \partial D;$$

(iv) if the Laplace transform M_3 exists, it is situated on the straight line $\{M_2, N\}$ with $N \in \{m, M_{-1} = \bar{M}_1, M_{-2} = \bar{M}_2\}$; (v) the tangent space of $M_2(p)$ is 1-dimensional for each $p \in \partial D$. Then $m(D)$ is a part of the Veronese surface.

Proof. Let us formulate our conditions analytically: (ii) means that F is an \mathbb{R} -valued function, i.e.,

$$(3.34) \quad F_2 = 0 \quad \text{on} \quad D;$$

(iv) gives

$$(3.35) \quad G = 0 \quad \text{on} \quad D;$$

(v) is equivalent to

$$(3.36) \quad E = B = C = 0 \quad \text{on} \quad \partial D;$$

for the last two conditions, see (3.23₂).

From (3.34) and (3.19), we get

$$(3.37) \quad dF_1 + F_1(\omega_0^0 - \omega_1^1) = 0.$$

The exterior differentiation yields, because of $F_1 \neq 0$,

$$(3.38) \quad A_1 H_2 + A_2 H_1 = 0.$$

Thus we have $AH = \bar{A}\bar{H}$, and (3.31), (2.25) imply

$$(3.39) \quad \psi_2 = \psi_{-2};$$

thus ψ_2 is a real form and

$$(3.40) \quad AH \neq 0.$$

From (3.29_{2,3}) and the condition (iii) we see that we may choose the frames in such a way that

$$(3.41) \quad F_1 = -1; \quad A_1 = \alpha = \frac{1}{3}\sqrt{3}, \quad A_2 = 0;$$

this and (3.38) imply

$$(3.42) \quad H_2 = 0, \quad H_1 \neq 0.$$

The condition (3.37) reduces then to

$$(3.43) \quad \omega_1^1 = \omega_0^0,$$

and (3.20) are simply

$$(3.44) \quad \begin{aligned} (\omega_3^3 - \omega_0^0) \wedge \omega^1 + (2\omega_1^2 + \omega_3^4) \wedge \omega^2 &= C_2 \omega^1 \wedge \omega^2, \\ (2\omega_1^2 + \omega_3^4) \wedge \omega^1 - (\omega_3^3 - \omega_0^0) \wedge \omega^2 &= C_1 \omega^1 \wedge \omega^2. \end{aligned}$$

Let functions $f_1, f_2: D \rightarrow \mathbb{R}$ satisfy a system of the form

$$(3.45) \quad \begin{aligned} df_1 \wedge \omega^1 + df_2 \wedge \omega^2 &= (a_{11}f_1 + a_{12}f_2) \omega^1 \wedge \omega^2, \\ df_2 \wedge \omega^1 - df_1 \wedge \omega^2 &= (a_{21}f_1 + a_{22}f_2) \omega^1 \wedge \omega^2, \end{aligned}$$

$a_{ij}: D \rightarrow \mathbb{R}$ being given. We may choose the coordinates (u, v) on D in such a way that

$$(3.46) \quad \omega^1 = r \, du, \quad \omega^2 = r \, dv; \quad r = r(u, v) \neq 0.$$

Then the system (3.45) may be rewritten as

$$(3.47) \quad f_{1v} - f_{2u} = -r(a_{11}f_1 + a_{12}f_2), \quad f_{1u} + f_{2v} = -r(a_{21}f_1 + a_{22}f_2).$$

This is clearly an elliptic system on D and $f_1 = f_2 = 0$ on ∂D implies $f_1 = f_2 = 0$ on D ; for this, see [9] or any other textbook on pseudoanalytic functions.

Applying this remark to (3.22_{1,2}) with $G_1 = G_2 = 0$, we see that $E_1 = E_2 = 0$ on D ; here we use (3.36₁). Similarly, we get $B_1 = B_2 = 0$ and $C_1 = C_2 = 0$. Thus (3.36) implies

$$(3.48) \quad E = B = C = 0 \quad \text{on } D.$$

Now the equations (3.21) are

$$(3.49) \quad \begin{aligned} \{dH_1 + H_1(\omega_0^0 - \omega_3^3)\} \wedge \omega^1 + (2\omega_1^2 + \omega_3^4) \wedge \omega^2 &= 0, \\ -(2\omega_1^2 + \omega_3^4) \wedge \omega^1 + \{dH_1 + H_1(\omega_0^0 - \omega_3^3)\} \wedge \omega^2 &= 0, \end{aligned}$$

and the system (3.19)–(3.22) reduces to (3.49) + (3.44) with $C_1 = C_2 = 0$. But this system immediately implies

$$(3.50) \quad dH_1 + (H_1 + 1)(\omega_0^0 - \omega_3^3) = 0,$$

this last equation being completely integrable. Applying Cartan's lemma to (3.44), we get the existence of functions M, N such that

$$(3.51) \quad \omega_3^3 - \omega_0^0 = M\omega^1 + N\omega^2, \quad 2\omega_1^2 + \omega_3^4 = N\omega^1 - M\omega^2,$$

with the differential consequences

$$(3.52) \quad \begin{aligned} (dM - N\omega_1^2) \wedge \omega^1 + (dN + M\omega_1^2) \wedge \omega^2 &= 0, \\ (dN + M\omega_1^2) \wedge \omega^1 - (dM - N\omega_1^2) \wedge \omega^2 &= -2(1 + 3\alpha H_1) \omega^1 \wedge \omega^2; \end{aligned}$$

for α , see (1.2). Thus we get functions K, L, P such that

$$(3.53) \quad dM - N\omega_1^2 = K\omega^1 + L\omega^2, \quad dN + M\omega_1^2 = L\omega^1 + P\omega^2,$$

$$(3.54) \quad K + P = 2(1 + 3\alpha H_1).$$

The exterior differentiation of (3.53) yields

$$(3.55) \quad \begin{aligned} (dK - 2L\omega_1^2) \wedge \omega^1 + \{dL + (K - P)\omega_1^2\} \wedge \omega^2 &= \\ = 2(1 + 2\alpha H_1) N\omega^1 \wedge \omega^2, \\ \{dL + (K - P)\omega_1^2\} \wedge \omega^1 + (dP + 2L\omega_1^2) \wedge \omega^2 &= \\ = -2(1 + 2\alpha H_1) M\omega^1 \wedge \omega^2, \end{aligned}$$

and we write

$$(3.56) \quad \begin{aligned} dK - 2L\omega_1^2 &= K_1\omega^1 + K_2\omega^2, \quad dP + 2L\omega_1^2 = P_1\omega^1 + P_2\omega^2, \\ dL + (K - P)\omega_1^2 &= L_1\omega^1 + L_2\omega^2 \end{aligned}$$

with

$$(3.57) \quad L_1 - K_2 = 2(1 + 2\alpha H_1)N, \quad P_1 - L_2 = -2(1 + 2\alpha H_1)M.$$

From (3.54) we get, using (3.50)

$$(3.58) \quad K_1 + P_1 = 6\alpha(H_1 + 1)M, \quad K_2 + P_2 = 6\alpha(H_1 + 1)N.$$

Consider the function $f: D \rightarrow \mathbb{R}$ defined by

$$(3.59) \quad 2f = M^2 + N^2.$$

Then

$$(3.60) \quad *df = -(ML + NP)\omega^1 + (MK + NL)\omega^2,$$

* being the Hodge *-operator with respect to the metric

$$(3.61) \quad ds^2 = (\omega^1)^2 + (\omega^2)^2 = -4\psi_1.$$

From (3.60) we get the Stokes theorem in the form

$$(3.62) \quad \int_{\partial D} *df = \int_D \{K^2 + 2L^2 + P^2 + 2(M^2 + N^2)(1 + 3\alpha + 5\alpha H_1)\} \omega^1 \wedge \omega^2.$$

Let us now calculate the Gauss curvature \varkappa' of the metric (3.61). We have

$$(3.63) \quad d\omega^1 = -\omega^2 \wedge \omega_1^2, \quad d\omega^2 = \omega^1 \wedge \omega_1^2, \quad d\omega_1^2 = -\varkappa' \omega^1 \wedge \omega^2,$$

i.e., $\varkappa' = 1 + 2\alpha H_1$. Thus the Gauss curvature \varkappa of $|\psi_1| = \frac{1}{4} ds^2$ is given by

$$(3.64) \quad \varkappa = 4(1 + 2\alpha H_1),$$

and we have

$$(3.65) \quad d\varkappa = 8\alpha(H_1 + 1)(M\omega^1 + N\omega^2)$$

because of (3.50) and (3.51). Because of (3.33₁) and (3.64),

$$(3.66) \quad 1 + 3\alpha + 5\alpha H_1 > 0 \quad \text{on } D.$$

The equation $H_1 = -1$ contradicts (3.66), and (3.33₂) implies $M = N = 0$ on ∂D .

Thus the integral formula (3.62) implies

$$(3.67) \quad M = N = 0 \quad \text{on } D.$$

The equations (3.51) reduce then to

$$(3.68) \quad \omega_3^3 - \omega_0^0 = 0, \quad 2\omega_1^2 + \omega_3^4 = 0$$

with the integrability condition $1 + 3\alpha H_1 = 0$, i.e.,

$$(3.69) \quad H_1 = -\alpha.$$

Finally, from (3.5), (3.17₁), (3.17₃) + (3.36₃), (3.43) and (3.68) we get

$$(3.70) \quad \omega_0^0 = \dots = \omega_4^4 = 0.$$

Considering now (3.15) + (3.17) with all the specializations made up to now, we see that we get exactly the equations of the form (1.3₂₋₆). QED.

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Author's address: Přehradní 10, 635 00 Brno, Czechoslovakia.